The gamma-ray burst variability–peak luminosity correlation: new results

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ABSTRACT

We report test results of the correlation between the time variability and the peak luminosity of gamma-ray bursts (GRBs), using a larger sample (32) of GRBs with known redshift than that available to Reichart et al., and using as a variability measure that introduced by these authors. The results are puzzling. Assuming an isotropic-equivalent peak luminosity, as had Reichart et al., a correlation is still found, but it is less relevant and inconsistent with a power law as previously reported. Assuming as the peak luminosity that corrected for GRB beaming for a subset of 16 GRBs with known beaming angle, the correlation becomes slightly less significant.

Key words: methods: data analysis – gamma-rays: bursts.

1 INTRODUCTION

Despite the small number of gamma-ray bursts (GRBs) with known redshift (several dozen), several correlations between intrinsic temporal or spectral parameters of the GRB prompt emission and GRB energetics have been discovered over the previous 7 years. Norris, Marani & Bonnell (2000) found an anticorrelation between the peak luminosity and the spectral lag (obtained by cross-correlating the time profiles of the same GRB in various energy bands), according to which more luminous bursts exhibit shorter lags. Salmonson & Galama (2002) discovered a positive correlation between the spectral lag of the gamma-ray prompt emission and the jet-break time of the afterglow decay, according to which a small break time corresponds to a small lag and consequently to a high peak luminosity of the GRB. Concerning the temporal properties of GRB time profiles, evidence has been found for a positive correlation between the temporal variability of the light curves and the isotropic-equivalent peak luminosity for GRBs with known redshift (Reichart et al. 2001, hereafter R01; Fenimore & Ramirez-Ruiz 2000).

Moreover, Reichart et al. (2003) have shown that the variability versus peak luminosity correlation could also hold true for X-ray flashes (XRFs; see Heise et al. 2001). As a consequence of the mentioned correlations, a correlation between time variability and spectral lag is also expected and confirmed for a large sample of Burst And Transient Source Experiment (BATSE) bursts (Schaefer, Deng & Band 2001). The variability versus peak luminosity correlation has been explained by several authors (e.g. Kobayashi, Ryde & MacFadyen 2002; Mészáros et al. 2002) mainly within the framework of the standard fireball model, according to which internal shocks between ultra-relativistic shells are responsible for the pulse-like structure of the GRB prompt emission, while the smooth afterglow emission is due to external shocks between the fireball wind and the matter surrounding the GRB progenitor (e.g. Piran 2004, for a review).

GRB variability-connected properties are thought to be more sensitive to the bulk Lorentz factor Γ and, if the GRB emission is beamed, to the jet opening angle and/or the viewing angle (e.g. Ioka & Nakamura 2001; Salmonson & Galama 2002). Within the fireball model, there are different mechanisms that could account for different time variabilities, also giving possible explanations for XRF properties (Mészáros et al. 2002). In addition, the ‘cannonball model’ for GRBs (Dado, Dar & De Rújula 2002) also seems to explain the variability versus peak luminosity correlation (Plaga 2001).

From the above correlations luminosity estimators have also been tentatively derived (Fenimore & Ramirez-Ruiz 2000; Reichart et al. 2001; Schaefer et al. 2001) to investigate general properties of GRBs, such as the luminosity function and the possible link with the star formation rate. In addition, empirical redshift indicators have been proposed based on the calibration derived with the small sample of GRBs with known redshift, making use of the X-ray and gamma-ray observations alone (Atteia 2003; Bagoly et al. 2003).

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In this work we test the variability versus peak luminosity correlation using the variability definition given by R01. We used a sample of 32 GRBs with known redshift. Furthermore, we studied the same correlation by replacing the isotropic-equivalent peak luminosity with that corrected for beaming for a subset of 16 GRBs with known collimation angle provided by Ghirlanda, Ghisellini & Lazatti (2004).

In Section 2, we discuss our sample of GRBs; in Section 3 we discuss the time variability analysis; in Section 4 we estimate the peak luminosity of the GRB in our sample and compare it with the R01 results. In Section 5 we present our results on variability–peak luminosity correlation and in Section 6 we discuss our results.

2 THE GRB SAMPLE

2.1 GRBs with known redshift

The sample of 32 GRBs with known redshift includes 16 GRBs detected by the gamma-ray burst monitor (GRBM) (Feroci et al. 1997; Froment et al. 1997; Costa et al. 1998) onboard the BeppoSAX satellite (Boella et al. 1997) during the period 1997–2002, two by the BATSE experiment (Paciesas et al. 1999) aboard the Compton Gamma–Ray Observatory (CGRO), six by the FREGATE instrument aboard HETE-II (Atteia et al. 2003), one by Konus/WIND (Aptekar et al. 1995), one by Ulysses (Hurley et al. 1997) and six by BAT/Swift (Gehrels et al. 2004). Eight of the 16 GRBs detected with the BeppoSAX GRBM were also detected with BATSE. We used public archives for GRB data obtained with BATSE, HETE-II, Konus/WIND and BAT/Swift. Table 1 reports the list of the GRBs in our sample including the spacecraft that detected them. When the same GRB has been detected by more than one instrument, we first checked the consistency of the results derived from different data sets and then concentrated on the instrument which had the best signal-to-noise ratio.

The time binning of the GRB light curves in our sample, which was used to derive the time variability, was the following: 7.8125 ms for the GRBM data, 64 ms for BATSE, 164 ms for HETE-II, 64 ms for Konus/WIND and 31.25 ms for Ulysses. In the case of BAT/Swift we made use of the event files and extracted the mask-tagged light curves with a binning time of 8 ms. Given that for the GRBs detected with the BeppoSAX GRBM, the high-resolution (7.8125-ms binning) time profiles are available in the 40–700 keV energy band, for the others we used the light curves in the energy bands which have the largest intersection with the 40–700 keV band: 110–320 keV (channel 3) for BATSE, 30–400 keV (band C) for FREGATE, 25–100 keV for Ulysses and 50–200 keV for Konus/WIND. In order to match the GRBM band, for BAT/Swift we extracted the light curves from the event files in the 40–350 keV band.

For GRB990510, given that the 7.8125-ms GRBM light curve is not available, we preferred to use 64-ms BATSE data rather than the 1-s GRBM light curve.

Six GRBs (980613, 011211, 021004, 050126, 050318 and 050416) with known redshift were not included in our sample due to their low signal, which prevented us from deriving a statistically significant variability estimate. Another GRB (021211, detected with HETE-II) was not included in the sample, due to the high ratio between the binning time and the smoothing time which could bias the variability estimate (see Section 3). Other GRBs with known redshift (000301C, 000418, 000926) detected by both Konus/WIND and Ulysses were not included in our sample, because unfortunately both Konus public and Ulysses data cover their light curves only partially.

In the case of BATSE (970828 and 001313) the usual four-channel 64-ms light curves are not available. Thus we made use of the 16-channel medium energy resolution (MER) spectra acquired along entire GRB with an integration time of 64 ms. Therefore, we rebinned the 16-energy-channel MER data both spectrally and temporally in order to reproduce as much as possible the four-channel scheme of BATSE 64-ms light profiles. As we discuss below, we relied on coarse time resolution light curves only when the overall duration of the GRB was very long compared with the binning time.

In general, the data available cover the entire light profile of the GRBs in our sample. However, there are some exceptions. In the case of the GRBM events, given that the high-resolution data cover 8 s before the trigger time and 98 s after it, in the case of the longest events (990506 and 010222, $T_{\text{90}} = 129$ and 97 s, respectively), it is not true. In these cases, the measure of time variability was obtained by summing the variability in the part covered by the 7.8125-ms bins with that in the part covered by 1-s ratemeters (the tail of the burst).

GRB000210 ($T_{\text{90}} = 8.1$ s) suffers from a 2.5-s long gap due to corrupted high-resolution data that occurred in the middle of the burst profile. Using the 1-s data, the mean 7.8125-ms light profile in the gap was reconstructed by adding Poisson noise to the mean profile. The value of GRB time variability thus derived was found not to change significantly, even when adding in the gap non-Poisson noise compatible with the 1-s profile.

For each GRB detected with the GRBM, we considered the light curves of the two most illuminated units and checked whether the best signal-to-noise ratio was obtained from a single unit or by summing the two units.

We found that for the 11 GRBs detected by GRBM and the Wide Field Cameras (WFCs; Jager et al. 1997), the best signal is obtained from a single GRBM unit (that with the larger area exposed to the GRB). For the five bursts detected with the GRBM but not with the WFCs, the best signal was obtained by summing the two most illuminated units: units 1 and 4 (980703), 3 and 4 (990506, 020405), and 2 and 4 (991216, 010921). In principle, the operation of adding the counts of different units is questionable because of dead time, as will be discussed in Section 3. In practice, for the above cases we made sure that the results were consistent with those obtained when considering only the most illuminated unit for each GRB. This has been found to be no longer true, i.e. the correction for dead time becomes non-negligible, when considering very small smoothing time-scales (Rossi et al., in preparation).

Concerning the eight bursts detected with both GRBM and BATSE (970508, 971214, 980425, 980703, 990123, 990506, 990510, 991216), we used the BeppoSAX data for 971214, 980703, 990123, 990506 and 991216, for which the higher time resolution of the GRBM turned out to be essential for a better variability estimate. While for the remaining three GRBs (970508, 980425 and 990510) we used the BATSE data given the better signal-to-noise ratio, after verifying the mutual consistency of the GRBM and BATSE variability results.

3 VARIABILITY MEASURE

We adopted the variability measure given by R01, slightly modified for two corrections which could affect the result: the instrument dead  

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1 ftp://cossc.gsfc.nasa.gov/compton/data/batse/ascii_data/64ms/
2 http://space.mit.edu/HETE/Bursts/Data/
3 http://heawww.gsfc.nasa.gov/docs/gamoscaray/legr/bacodine/konus_grbs.html
4 http://swift.gsfc.nasa.gov/
Table 1. Variability versus the peak luminosity for 32 GRBs with known redshift (1σ errors).

<table>
<thead>
<tr>
<th>GRB name</th>
<th>z redshift</th>
<th>Missiona</th>
<th>$T_{f=0.45}$ (s)</th>
<th>$V_{f=0.45}$</th>
<th>Peak lum. $L^b$ (10$^{50}$ erg s$^{-1}$)</th>
<th>z Refs.c</th>
</tr>
</thead>
<tbody>
<tr>
<td>970228</td>
<td>0.695</td>
<td>BS/U/K</td>
<td>2.2</td>
<td>0.223±0.018</td>
<td>48.7±9.9</td>
<td>1</td>
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<tr>
<td>970508</td>
<td>0.835</td>
<td>B/BS/U/K</td>
<td>2.4</td>
<td>0.023±0.013</td>
<td>9.43±1.89</td>
<td>2</td>
</tr>
<tr>
<td>970828</td>
<td>0.958</td>
<td>B/BS/U/K</td>
<td>12.9</td>
<td>0.101±0.002</td>
<td>120.0±40.0</td>
<td>3</td>
</tr>
<tr>
<td>971214</td>
<td>3.418</td>
<td>BS/B/UK/N/R</td>
<td>4.4</td>
<td>0.110±0.012</td>
<td>360.0±65.0</td>
<td>4</td>
</tr>
<tr>
<td>980425</td>
<td>0.0085</td>
<td>B/BS/U/K</td>
<td>4.7</td>
<td>0.043±0.048</td>
<td>0.0007±0.0002</td>
<td>5</td>
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<tr>
<td>980703</td>
<td>0.966</td>
<td>B/BS/U/K</td>
<td>3.2</td>
<td>0.044±0.007</td>
<td>26.4±5.6</td>
<td>6</td>
</tr>
<tr>
<td>990123</td>
<td>1.6</td>
<td>BS/B/UK</td>
<td>12.8</td>
<td>0.112±0.002</td>
<td>840.0±121.0</td>
<td>7</td>
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<tr>
<td>990506</td>
<td>1.3</td>
<td>BS/B/UK/R</td>
<td>8.6</td>
<td>0.270±0.005</td>
<td>583.0±121.0</td>
<td>8</td>
</tr>
<tr>
<td>990510</td>
<td>1.619</td>
<td>BS/B/UK/N</td>
<td>3.2</td>
<td>0.214±0.005</td>
<td>300.0±50.0</td>
<td>9</td>
</tr>
<tr>
<td>990705</td>
<td>0.86</td>
<td>BS/U/UK/N</td>
<td>8.0</td>
<td>0.178±0.003</td>
<td>134.0±21.0</td>
<td>10, 11</td>
</tr>
<tr>
<td>990712</td>
<td>0.434</td>
<td>BS/U/K</td>
<td>4.1</td>
<td>0.04±0.017</td>
<td>5.4±1.0</td>
<td>12</td>
</tr>
<tr>
<td>991208</td>
<td>0.706</td>
<td>K/U/N</td>
<td>5.1</td>
<td>0.08±0.003</td>
<td>290.0±100.0</td>
<td>13</td>
</tr>
<tr>
<td>991216</td>
<td>1.02</td>
<td>BS/B/UK/N</td>
<td>2.6</td>
<td>0.193±0.002</td>
<td>1398.0±200.0</td>
<td>14</td>
</tr>
<tr>
<td>990131</td>
<td>4.5</td>
<td>B/BS/U/K</td>
<td>8.0</td>
<td>0.18±0.005</td>
<td>3600.0±900.0</td>
<td>15</td>
</tr>
<tr>
<td>990210</td>
<td>0.846</td>
<td>BS/U/K</td>
<td>1.59</td>
<td>0.026±0.002</td>
<td>80.0±50.0</td>
<td>16</td>
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<tr>
<td>990911</td>
<td>1.058</td>
<td>U/K/N</td>
<td>5.2</td>
<td>0.07±0.034</td>
<td>360.0±60.0</td>
<td>17</td>
</tr>
<tr>
<td>990222</td>
<td>1.477</td>
<td>BS/U/K</td>
<td>6.62</td>
<td>0.201±0.003</td>
<td>801.0±119.0</td>
<td>18</td>
</tr>
<tr>
<td>990921</td>
<td>0.42</td>
<td>BS/H/UK</td>
<td>5.3</td>
<td>0.038±0.016</td>
<td>8.0±2.0</td>
<td>19</td>
</tr>
<tr>
<td>991121</td>
<td>0.36</td>
<td>BS/U/UK/O</td>
<td>8.3</td>
<td>0.049±0.002</td>
<td>19.9±3.1</td>
<td>20</td>
</tr>
<tr>
<td>990214</td>
<td>3.198</td>
<td>H/UK</td>
<td>8.8</td>
<td>0.203±0.017</td>
<td>300.0±60.0</td>
<td>21</td>
</tr>
<tr>
<td>990405</td>
<td>0.69</td>
<td>BS/U/UK/O</td>
<td>9.9</td>
<td>0.169±0.007</td>
<td>71.4±11.2</td>
<td>22</td>
</tr>
<tr>
<td>990813</td>
<td>1.25</td>
<td>H/UK/O</td>
<td>17.4</td>
<td>0.248±0.007</td>
<td>340.0±70.0</td>
<td>23</td>
</tr>
<tr>
<td>990326</td>
<td>1.98</td>
<td>H/UK</td>
<td>26.6</td>
<td>0.045±0.015</td>
<td>25.0±5.0</td>
<td>24</td>
</tr>
<tr>
<td>990328</td>
<td>1.52</td>
<td>H/UK</td>
<td>24.9</td>
<td>0.051±0.005</td>
<td>90.0±18.0</td>
<td>25</td>
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<tr>
<td>990329</td>
<td>0.168</td>
<td>H/UK/OK/R</td>
<td>4.9</td>
<td>0.105±0.007</td>
<td>6.1±1.2</td>
<td>26</td>
</tr>
<tr>
<td>990406</td>
<td>0.712</td>
<td>H/R/HH</td>
<td>8.0</td>
<td>0.053±0.002</td>
<td>66.0±10.0</td>
<td>27</td>
</tr>
<tr>
<td>990315</td>
<td>1.949</td>
<td>BSw</td>
<td>12.3</td>
<td>0.04±0.032</td>
<td>38.0±8.0</td>
<td>28</td>
</tr>
<tr>
<td>990319</td>
<td>3.24</td>
<td>BSw</td>
<td>3.6</td>
<td>0.06±0.032</td>
<td>84.0±20.0</td>
<td>29</td>
</tr>
<tr>
<td>990401</td>
<td>2.90</td>
<td>BSw</td>
<td>4.4</td>
<td>0.19±0.028</td>
<td>740.0±100.0</td>
<td>30</td>
</tr>
<tr>
<td>990505</td>
<td>4.27</td>
<td>BSw</td>
<td>9.0</td>
<td>0.20±0.043</td>
<td>250.0±50.0</td>
<td>31</td>
</tr>
<tr>
<td>990525</td>
<td>0.606</td>
<td>BSw</td>
<td>2.0</td>
<td>0.11±0.003</td>
<td>8.0±10.0</td>
<td>32</td>
</tr>
<tr>
<td>990603</td>
<td>2.821</td>
<td>BSw</td>
<td>1.2</td>
<td>0.24±0.034</td>
<td>1200.0±200.0</td>
<td>33</td>
</tr>
</tbody>
</table>

Notes. aMission: BS (BeppoSAX), B (BATSE/CGRO), K (Konus/WIND), H (HETE-II), U (Ulysses), S (SROSS-C), N (NEAR), R (RossiXTE), O (Mars Odyssey), RH (RHESSI), BSw (BAT/Swift); the data used are taken from the first mission mentioned.
bIsotropi-equivalent peak luminosity in 10$^{50}$ erg s$^{-1}$ in the rest frame 100–1000 keV band, for peak fluxes measured on a 1-s time-scale, $H_0 = 65$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_{m} = 0.3$ and $\Omega_{\Lambda} = 0.7$.

time and a small amount of non-Poisson noise present in the GRBM background data. The variability measure used by R01 was defined as a properly normalized mean square deviation of the intrinsic light curve of a GRB in a given energy band from a smoothed one. For a discrete light curve formed from $N$ bins, the variability measure, according to R01, is given by

$$V_{f,p} = \frac{\sum_{i=1}^{N}[S(C_i, N_i) - S(C_i, N_p)]^2}{\sum_{i=1}^{N}[S(C_i, N_i) - B_i]^2}$$

(1)
where by the intrinsic light curve we mean the GRB light curve in the source frame, \( N_j \) is the number of data bins corresponding to the smoothing time-scale \( T_j \) defined by R01 as the shortest cumulative time interval during which a fraction \( f \) of the total counts above background has been collected, \( C_j \) and \( B \) are the original GRB (source plus background) and background counts in bins \( j \) and \( i \), respectively, in the observer frame, the index \( P \) means that the variability measure is inclusive of the Poisson noise. \( S_i(C_j, N_i) \) is roughly the mean counts on \( N_i \) bins (\(x = z \) or \( f \)) centred around the \( i \)th bin:

\[
S_i(C_j, N_i) = \frac{1}{N_i} \sum_{j-i}^{j+i} C_j + \left( \frac{N_i - 1}{2} - n_i \right) \left( C_{i-n_i+1} + C_{i+n_i+1} \right). \tag{2}
\]

\( N_i \) is the number of bins in the observer frame, which corresponds to one bin in the source frame. Assuming as the time duration of one bin in the source frame the shortest binning \( \Delta t \) of the data (e.g. in the case of the BeppoSAX GRBM, \( \Delta t = 7.8125 \) ms), in the observer frame the number of bins, depending on the GRB redshift \( z \), for relativistic time dilation and narrowing of the light curves at high energies (Fenimore et al. 1995), is given by \( N_i = (1+z)^{\delta} \) with \( \beta \approx 0.6 \). Thus \( N_i \) can take non-integer values and \( n_i \) is the truncated integer value of \((N_i - 1)/2\).

R01 found that the best luminosity estimator is obtained when using \( f = 0.45 \); for this reason, we fixed \( f = 0.45 \).

The variability \( V_{f,P} \) can also be written as follows:

\[
V_{f,P} = \frac{\sum_{i=1}^{N} \left( \frac{\sum_{j=1}^{N} a_{ij} C_j}{\sum_{j=1}^{N} b_{ij} C_j - B_j} \right)^2}{\sum_{i=1}^{N} \left( \frac{\sum_{j=1}^{N} b_{ij} C_j - B_j}{\sum_{j=1}^{N} b_{ij} C_j} \right)^2}.
\tag{3}
\]

where the coefficients \( a_{ij} \) and \( b_{ij} \), for each GRB, are computed by comparing equation (3) with equation (1) through equation (2).

Following R01, after subtraction of the Poisson variance the variability measure is given by

\[
V_f = \frac{\sum_{i=1}^{N} \left[ \left( \frac{\sum_{j=1}^{N} a_{ij} C_j}{\sum_{j=1}^{N} b_{ij} C_j - B_j} \right)^2 - \sum_{j=1}^{N} a_{ij}^2 C_j \right]}{\sum_{i=1}^{N} \left[ \frac{\sum_{j=1}^{N} b_{ij} C_j - B_j}{\sum_{j=1}^{N} b_{ij} C_j} \right]^2 - \sum_{j=1}^{N} b_{ij}^2 C_j}. \tag{4}
\]

which is the expression used by R01 to evaluate the variability of the GRBs in their sample. We slightly modified the above expression by also taking into account the dead time, which is known to affect the Poisson variance of a stationary process (Müller 1973, 1974; Libert 1978). In the case of a stationary Poisson process with measured mean rate \( \mu \), the variance of its counts in the time bin \( \Delta t \), which is given by \( \mu \Delta t \) in the absence of dead time, becomes \( \mu \Delta t (1 - \mu \Delta t)^2 \) in the asymptotic limit \( \tau/\Delta t < 1 \), where \( \tau \) is the dead time. In the case of the BeppoSAX GRBM \( \tau = 4 \) ms, \( \tau/\Delta t \approx 5 \times 10^{-4} \) for the shortest bin duration \( \Delta t = 7.8125 \) ms. In the same limit \( \tau/\Delta t \ll 1 \), the same correction factor \( (1 - \mu \Delta t)^2 \) applies to the white noise level of the power spectral density (PSD) estimate (Frontera & Fuligni 1978; van der Klis 1989). It is shown (Frontera & Fuligni 1979) that the same correction factor holds when the process is non-stationary, such as GRBs or flares. Potentially, our variability calculations could be sensitive to dead time, especially for those GRBs with huge peak count rates, such as in the case of 990123 (\( \sim 16\,000 \) count s\(^{-1} \) with GRBM), for which, around the peak, the true variance is \( \sim 0.9 \) times the measured counts.

In addition, we corrected for a slight (a few per cent) degree of non-Poisson noise found in the GRBM high-resolution data. This noise increases the Poisson variance by a factor \( r_{np} \) which ranges from 1.027 to 1.049, depending on the detection unit, for the GRBM data after 1996 November.\(^3\)

Taking into account both dead time and non-Poisson noise, the right terms to be subtracted in the numerator and denominator of equation (3) become \( \sum_{j=1}^{N} a_{ij}^2 C_j r_j \) and \( \sum_{j=1}^{N} b_{ij}^2 C_j r_j \) respectively (see, for comparison, equation 4), where

\[
r_j = r_{np} \left( 1 - C_j \frac{\tau}{\Delta t} \right)^2. \tag{5}
\]

Consequently, the expression we used to estimate the net GRB time variability is given by

\[
V_f = \frac{\sum_{i=1}^{N} \left[ \left( \frac{\sum_{j=1}^{N} a_{ij} C_j}{\sum_{j=1}^{N} b_{ij} C_j - B_j} \right)^2 - \sum_{j=1}^{N} a_{ij}^2 C_j r_j \right]}{\sum_{i=1}^{N} \left[ \frac{\sum_{j=1}^{N} b_{ij} C_j - B_j}{\sum_{j=1}^{N} b_{ij} C_j} \right]^2 - \sum_{j=1}^{N} b_{ij}^2 C_j r_j}. \tag{6}
\]

We used as the statistical uncertainty \( \sigma_{V_f} \) on the variability measure given by R01 (equation 8) properly modified to take into account the correction factor \( r_j \).

We found that the variability measure is not sensitive to dead time corrections for long GRBs, in which \( T_f > 0.45 \) is much longer than the bin time, while it is significantly modified for relatively short GRBs exhibiting sharp intense pulses.

### 3.1 Variability dependence on binning time

In order to understand how time binning affects the GRB variability, for the brightest GRBs, detected with either GRBM or BATSE we evaluated the variability measure (equation 6) as a function of the binning time of the data. The result is that the variability is better estimated for very short time durations of the data bins with respect to the smoothing time-scale \( T_s = T_{f > 0.45} \). More specifically, it results that the variability significantly decreases for few 0.01 < \( \Delta t/T_s < 0.1 \), and becomes unreliable when this ratio becomes still higher, i.e., for \( \Delta t/T_s > 0.1 \).

On the other hand, the bin time should be long enough to collect a significant number of photons (typically at least 20 bin\(^{-1}\) on average) to ensure the Gaussian limit and take account of the effects of statistical fluctuations. Thus we rejected those GRBs where the data sets do not match the above requirements.

Figs 1 and 2 show the illustrative cases of 991216 and 970228, respectively. We calculated \( V_{f > 0.45} \) using both GRBM and BATSE (KONUS) data sets as a function of the binning time for 991216 (970228). For both GRBs, the variability seems to approach an asymptotic value for decreasing values of binning time. In the case of 991216, it appears that the original binning time of BATSE, 64 ms, is a little too coarse as its corresponding value of \( V_{f > 0.45} \) is significantly lower: we assume as an asymptotic value the measure obtained with the smallest binning time of GRBM data and obtain \( V_{f > 0.45} = 0.193 \pm 0.002 \), while the BATSE measure is 0.170 ± 0.003, i.e., ~6σ apart. In contrast, in the case of 970228 the KONUS measure with the smallest binning time of 64 ms yields a measure of \( V_{f > 0.45} \) which is apparently consistent with the GRBM one.

\(^3\) During the first months of BeppoSAX operation the non-Poisson noise of the GRBM was much higher, due to the too low energy threshold (around 20 keV) set at the beginning of the mission (Feroci et al. 1997).
GRB variability–peak luminosity correlation

4 PEAK LUMINOSITY ESTIMATE

The GRB peak luminosities were estimated using the definition of luminosity distance in the source frame 100–1000 keV energy band:

\[ L = 4\pi D_L^2(z) \int_{100}^{1000} E \Phi(E) \, dE, \]

(7)

where \( \Phi(E) \) is the measured spectrum (photon cm\(^{-2}\) s\(^{-1}\) keV\(^{-1}\)) around the peak time, \( D_L(z) \) is the luminosity distance at redshift \( z \) and \( E \) is the energy expressed in keV. By replacing \( E' = E(1 + z) \)

we obtain

\[ L = \frac{4\pi D_L^2(z)}{(1+z)^2} \int_{100}^{1000} E' \Phi \left( \frac{E'}{1+z} \right) \, dE'. \]

(8)

Formally, equation (8) is the same as equation (9) of R01: there \( D(z) \) is the comoving distance, which is equal to \( D_L(z) = D_L(1 + z) \) if we consider a flat universe. However, unlike R01 who used as \( \Phi(E) \) the best-fitting Band model (Band et al. 1993) to the average GRB count spectrum normalized to the peak count rate, we used for the GRBM data the best-fitting power-law spectrum \[ \Phi(E) = NE^{-\alpha} \]

to the GRB peak count rate spectrum obtained from the 1-s ratemeters available in two channels (40–700 and \( >1000 \) keV). When the 225-channel time-averaged spectrum was not available, we added a conservative 10 per cent systematic error to the peak luminosity uncertainties. Thus, for the GRBM bursts, equation (8) becomes

\[ L = 4\pi D_L^2(z)(1+z)^{\alpha-2} F_p, \]

(9)

where \( F_p = \int_{100}^{1000} NE^{-\alpha} \, dE' \) is the 100–1000 keV peak flux measured in the observer frame (erg cm\(^{-2}\) s\(^{-1}\)). In the case of GRBs with sharp peaks of \( <1 \) s duration (e.g. GRB000214), the peak luminosity obtained from 1-s ratemeters was further corrected by the ratio between the actual peak value and that derived from 1-s ratemeters.

For the GRBs in our sample not detected with the GRBM, we used the best-fitting parameters of \( \Phi(E) \) available from the literature. The best-fitting spectral parameters for \textit{HETE}-II bursts were taken from Sakamoto et al. (2005), except for the recent 041006 for which we used the best-fitting cut-off power-law parameters published by the \textit{HETE}-II team on the HETE web page \( (E_0 = 100.2 \text{ keV}, \alpha = 1.367) \). For the \textit{Ulysses} GRB000911 we made use of the best-fitting parameters published by Price et al. (2002a), while for the Konus burst 991208 we used the parameter values given by R01. For the \textit{BATSE} GRB000131 we fitted the peak energy spectrum from MER data in the range 30–1000 keV with the Band function \( (\alpha = -0.56, \beta = -2.17, E_0 = 153 \text{ keV}, \chi^2/\text{dof} = 1.0) \). Likewise, for the \textit{BATSE} GRB970828 the peak energy spectrum was fitted with the Band function \( (\alpha = -0.65, \beta = -2.56, E_0 = 269 \text{ keV}, \chi^2/\text{dof} = 1.1) \). For \textit{BAT/Swift} we extracted from the event file the 1-s 80-channel spectrum around the peak; for all six \textit{BAT/Swift} GRBs considered, respectively. In the case of the couple of GRBs considered above, we infer that GRBM data turned out to be essential in estimating the variability of 991216, as \textit{BATSE} data alone, although consistent with GRBM data for comparable binning times, do not seem to approach an asymptotic value of \( V_{f=0.45} \), while GRBM data do.

On the other hand, in the case of 970228 KONUS data exhibit an asymptotic trend towards small binning times; together with the fulfilment of the other two requirements, KONUS time resolution is acceptable and yields a variability measure which is consistent with the GRBM within errors.
the peak spectrum was fitted with a simple power law in the 15–350 keV range, apart from a couple of them (050525 and 050603) for which only the cut-off power-law model yields a good fit. We then used equation (7) to evaluate the peak luminosity.

Our peak luminosity estimates are reported in Table 1.

For the common sample of GRBs, our estimates of the peak luminosity are fully consistent with those obtained by R01 (see Fig. 3).

5 RESULTS

5.1 GRBs with known redshift

First of all, we evaluated the time variability of the GRBs (13) common to our sample and to that by R01, in order to test the mutual consistency of our results with those obtained by R01.

5.1.1 Variability

Fig. 4 compares the two time variability estimates. As can be seen, the results are well consistent with each other, except for three cases (970228, 991216 and 000131).

For each of these GRBs we investigated the reason for the discrepant measure of $V_{f=0.45}$: first of all by trying to reproduce the results of R01 using Konus data alone for 970228 and BATSE data alone for the other two.

5.1.2 GRB 970228

In order to reproduce R01’s results for 970228 we used the same data set, i.e. the light curve by Konus. The only difference is that we used public data that include a single light curve in the 50–200 keV energy band, while R01 used three different energy bands: 10–45, 45–190 and 190–770 keV. R01 report the smoothing time-scale for each of the three energy channels and only the global variability measure derived from merging the three different measures according to the procedure described therein. Our measure of the smoothing time-scale is $2.82 \pm 0.32$ s to be compared with that obtained by R01 for the same energy channel, i.e. 2.891 s (no error is reported), and thus is consistent. Our variability measure with Konus data is $0.19 \pm 0.04$ to be compared with R01’s one of $0.08 \pm 0.05$. The measure obtained with GRBM data, $0.22 \pm 0.02$, agrees well with our Konus measure (see Fig. 2), but does not with the R01 Konus one. The measure reported by R01 was derived from the three energy channels; this might partially explain the difference. However, we note that our Konus measure is 2.2σ apart from the R01 value of 0.08. We are led to think of two potential sources of discrepancy between our measure and R01’s. First, the overall time interval containing the GRB might be different; secondly, the extrinsic scatter that R01 find on the variability measure. We refer the reader to the R01 paper for a definition of the extrinsic scatter of variability due to the different energy channels derived for each GRB. As we neither have the same Konus data as R01, nor do we know the overall time interval adopted by R01, we cannot establish conclusively the reason for the discrepancy for this GRB. However, concerning the first possibility, we tentatively adopted other time intervals trying to match the variability measure reported by R01. We find a variability measure of $0.08 \pm 0.03$ for a time interval including the first sharp pulse and lasting $\sim 40$ s until the first pulse following a quiescent interval from the very first pulse. It must be pointed out that our true measure was performed on a 80-s long interval, as there is evidence for emission. This could be a hint for the possible explanation of the discrepancy.

5.1.3 GRB 991216

For this GRB we adopted the measure obtained with GRBM data and we have already discussed the reasons for this in Section 3.1. Here we try to reproduce the R01 results using the same BATSE data and then compare our variability measures on each energy channel with the merged value derived by R01. In Fig. 5 we show the variability as a function of the BATSE energy channel and compare them with the
merged value \( \pm 1\sigma \) reported by R01. We remind the reader that when comparing with GRBM results, we just considered channel 3. In particular, for this GRB, we know from previous discussion that the value obtained with BATSE channel 3 appears to be underestimated with respect to the GRBM result (see Fig. 1).

We have a perfect match within errors between our set of four smoothing time-scale values and those obtained by R01. Therefore, we are led to conclude that our variability measures should match consequently. On this basis, from Fig. 5 we note that! the extrinsic scatter by R01, whose \( 1\sigma \) region is displayed through dashed lines, seems to be slightly underestimated. In fact, the channel 1 measure is \( 2.6\sigma \) below and the channel 3 is \( 3\sigma \) above. In addition to this, we remind the reader that for this particular GRB exhibiting sharp pulses we know from GRBM data that a time binning of 64 ms is too coarse (see Fig. 1 and the discussion in Section 3.1). We therefore conclude that the effect of a higher scatter of variability at different energy channels than that estimated by R01, combined with the fact that for this GRB a binning time of 64 ms seems inadequate, account for the discrepancy between our measure of variability for 991216 and that published by R01.

5.1.4 GRB 000131

For this GRB we made use of BATSE data while R01 used Konus data. Unfortunately, the public Konus data of this GRB do not cover the whole profile, so we are bound to use BATSE data alone and compare our variability measures with the R01 value. Fig. 6 displays the variability as a function of the BATSE energy channel and dashed lines show the R01 estimate. The reasons for the discrepancy, which is apparent from Fig. 6, are due to the different smoothing time-scales: by comparing our set of four values with the three corresponding to the three lower Konus channels (the light curve of channel 4 cannot be used according to R01), our values are systematically greater than R01’s. If we adopt the same time-scales obtained by R01 we obtain variability measures which are consistent with the R01 value. This conclusively proves that the discrepancy for this GRB must be ascribed to the different measures of the time-scales. Concerning the origin of this discrepancy in the time-scale evaluation, we do not find any apparent bias that could have affected the calculations using BATSE profiles.

5.1.5 GRB 050315: a BAT/Swift GRB

The variability measured for this BAT/Swift GRB is consistent with that published by Donaghy et al. (2005). Here we want to show the consistency of the variability measure with BAT/Swift data and its dependence on the different BAT energy channels. Fig. 7 shows the variability as a function of the BAT channels obtained by us (asterisks and solid lines) and by Donaghy et al. (2005, squares and dashed lines). The energy bands of the three BAT channels considered are the following: 15–25, 25–50 and 50–100 keV, respectively. Channel 5 in Fig. 7 corresponds to the integrated band 15–100 keV. These energy channels have been chosen in order to match those used by Donaghy et al. (2005). Clearly, the two sets of variability measures are consistent within errors for each single BAT channel.

5.1.6 Variability–peak luminosity

Fig. 8 and Table 1 show the \( V_{f=0.45} \) versus peak luminosity for the entire sample of 32 GRBs with known redshift. Dashed lines...
of the best-fitting power-law relationship found by R01 along with the ±1σ width, according to which $L \propto V_{\text{rf}}^{m}$, where $m = 3.3^{+1.1}_{-0.9}$. Apparently, from Fig. 8, the correlation between the GRB variability and the peak luminosity is confirmed, as also demonstrated by the correlation coefficients and their statistical significance. The results are given in Table 2, where we report the values of both the Pearson linear correlation coefficient $r$, and the non-parametric correlation coefficients $r_s$ (the Spearman rank-order coefficient) and $\tau$ (the Kendall coefficient; Press et al. 1993), along with the corresponding correlation statistical significance. The same correlation coefficients have been evaluated also taking into account error bars on both $V_{\text{rf}}$ and $L$ through simulations (reported in parentheses in Table 2). We scattered each point along with its error bars assuming a Gaussian probability distribution in both dimensions and then we calculated the mode for each coefficient distribution.

However, the high spread of the data points, clustered in two main regions of the parameter space, shows that not only are the best-fitting power-law parameters obtained by R01 in disagreement with the our results but also that a power-law model gives a poor description of the data. Indeed, by fitting the data with a power-law model:

$$\log L_{50} = m \log V_{\text{rf}} + q,$$

where the peak luminosity is $L = L_{50} \times 10^{50}$ erg cm$^{-2}$ s$^{-1}$, independently of the method used for the fit (the usual least-squares

Table 2. Correlation coefficients for GRBs with known redshift. We also report in parentheses the mode values obtained from simulations.

<table>
<thead>
<tr>
<th>Kind</th>
<th>Coefficient</th>
<th>Probability</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pearson's $r$</td>
<td>0.514 (0.412)</td>
<td>0.0026 (0.019)</td>
</tr>
<tr>
<td>Spearman's $r_s$</td>
<td>0.625 (0.612)</td>
<td>0.0001 (0.0002)</td>
</tr>
<tr>
<td>Kendall’s $\tau$</td>
<td>0.446 (0.436)</td>
<td>0.0003 (0.0005)</td>
</tr>
</tbody>
</table>

Table 3. Best-fitting power-law parameters of the $L$ versus the $V_{\text{rf}}$ correlation for GRBs with known redshift.

<table>
<thead>
<tr>
<th>Method</th>
<th>$m$</th>
<th>$q$</th>
<th>$\chi^2$/dof</th>
</tr>
</thead>
<tbody>
<tr>
<td>Least-squares fit</td>
<td>$1.30^{+0.84}_{-0.44}$</td>
<td>$3.36^{+0.89}_{-0.43}$</td>
<td>1167/30</td>
</tr>
<tr>
<td>Least-absolute-deviation fit</td>
<td>$1.16^{+0.53}_{-0.17}$</td>
<td>$3.32^{+0.49}_{-0.15}$</td>
<td>1145/30</td>
</tr>
</tbody>
</table>

Figure 8. $V_{\text{rf}} = 0.45$ versus the peak luminosity for GRBs with known redshift. Dashed lines mark the best-fitting power-law relationship found by R01 (central line) and ±1σ widths.

Figure 9. Contour plot of the 1σ region of the two best-fitting parameters $m$ and $q$ (least-squares fit) in the case of $V_{\text{rf}} = 0.45 – L$. 

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Figure 10. Top panel, beaming-corrected rest frame peak luminosity $L_{p,\gamma}$ versus the variability for a subset of 16 GRBs with known redshift and beaming angle (Ghirlanda et al. 2004). Also shown are two lower limits (971214 and 011121) and three upper limits (000131, 000911 and 010921). Bottom panel, $L_p$ isotropic-equivalent peak luminosity versus the variability for the same 16 GRBs.

fit or minimization of absolute deviations, see Press et al. 1993), we find unsatisfactory results ($\chi^2 > 1000$, 30 dof, in either case). Just to compare our best-fitting power-law model results with those obtained by R01, in Table 3 we report the best-fitting parameters of the power law for the two above-mentioned fitting methods.

In Fig. 9 we report the $1\sigma$ contour plot of the best-fitting parameters $m$ and $q$. As can be seen, the two parameters are highly correlated.

We also evaluated the statistical uncertainty in $\log L_{50}$ as a function of $V_f=0.45$, taking into account the correlation between the two
parameters. In Fig. 8 the point corresponding to GRB 980425 is out of the plot window to avoid scale compression, but its variability is affected by a large uncertainty (0.049–0.048, see Table 1).

5.2 Luminosity correction for GRB beaming

For the GRBs with known redshift, we also investigated the correlation between variability $V_{f=0.45}$ and peak luminosity, after correcting the luminosity values given in Table 1 for the GRB beaming angles estimated by Ghirlanda et al. (2004). Ghirlanda et al. (2004) demonstrated that after this correction, the correlation between the peak energy $E_{p}^{\text{iso}}$ in the source frame and $E_{p,\text{c}}$ (corrected for beaming) improved with a lower spread of the data point around the best-fitting curve. From our sample of GRBs with known redshift, we considered those for which Ghirlanda et al. (2004) provided the beaming angles: the resulting subset includes 16 GRBs. In our case the result is shown in Fig. 10. Unlike the findings by Ghirlanda et al. (2004) for the Amati et al. (2002) relation, in the present case the spread of the data points becomes larger when the energy released in the GRB is corrected for beaming, although the correlation remains significant to within 1 per cent. We computed the correlation coefficients for this subset of GRBs in both cases: either assuming beam-corrected $L_{p,\gamma}$ and isotropic equivalent $L_{p}$ peak luminosities versus variability (see Table 4).

As reported in Table 4 and clearly shown by Fig. 10, in the case of the isotropic-equivalent peak luminosity the spread is smaller than in the case when the correction for beaming is applied.

6 DISCUSSION

The results found are puzzling. We confirm the peak luminosity versus variability correlation found by R01 when we use their sample of GRBs, but we find a much larger spread of the data points when a larger sample (32 events) of GRBs is used. In this case the correlation between $V_{f=0.45}$ and $L_{\text{iso}}$ is confirmed (at a significance of $\lesssim 3 \times 10^{-4}$ according to non-parametric tests), but the data points are spread out in only two regions of the parameter space, with a poor description of the data points ($\chi^{2} > 1000$, 30 dof) with a power-law function, which was the best-fitting function found by R01. If, in spite of that, this function is used as a fitting model, the power-law index derived from our data ($\gamma = 0.048$) cannot be described by a power-law function as found by R01. If, in spite of that, we fit the data with this function we find that the power-law index ($\gamma = 0.048$) is much lower and inconsistent with that found by R01 ($\gamma = 0.34$). If we correct the peak luminosity for the GRB beaming, the correlation is less significant.

7 CONCLUSIONS

We have tested the correlation found by R01 between peak luminosity and time variability following the same method used by R01 with a larger sample of GRBs. For 32 GRBs with known redshift we confirm the existence of a correlation between the measure of time variability defined by R01 and the isotropic-equivalent peak luminosity. However, we find a much higher spread of the data points in the parameter space, with the consequence that the correlation cannot be described by a power-law function as found by R01. If, in spite of that, we fit the data with this function we find that the power-law index ($\gamma = 0.048$) is much lower and inconsistent with that found by R01 ($\gamma = 0.34$). If we correct the peak luminosity for the GRB beaming, the correlation is less significant.

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For a couple of them, i.e. 041006 and 050525, the values derived by the same authors are taken from the following web site: http://www.merate.mi.astro.it/~ghirla/deep/blink.htm

Table 4. Correlation coefficients for 16 GRBs with known redshift and beaming angle: $V$ versus the beaming-corrected $L_{p,\gamma}$ (first two columns) and $V$ versus the isotropic-equivalent $L_{p}$ (last two columns).

<table>
<thead>
<tr>
<th>Kind</th>
<th>$V$ versus $L_{p,\gamma}$</th>
<th>$V$ versus $L_{p}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pearson’s $r$</td>
<td>0.664</td>
<td>0.005</td>
</tr>
<tr>
<td>Spearman’s $r_{s}$</td>
<td>0.653</td>
<td>0.006</td>
</tr>
<tr>
<td>Kendall’s $t$</td>
<td>0.467</td>
<td>0.012</td>
</tr>
</tbody>
</table>

$\gamma$ for a couple of them, i.e. 041006 and 050525, the values derived by the same authors are taken from the following web site: http://www.merate.mi.astro.it/~ghirla/deep/blink.htm
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