Blend lines in the polarized spectrum of the Sun

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ABSTRACT

Blend lines form an integral part of the theoretical analysis and modelling of the polarized spectrum of the Sun. Their interaction with other spectral lines needs to be explored and understood before we can properly use the main spectral lines to diagnose the Sun. They are known to cause a decrease in the polarization in the wings of the main line on which they superpose, or in the polarization of the continuum, when they are assumed to be formed either under the local thermodynamic equilibrium (LTE) conditions or when their intrinsic polarizability factor is zero. In this paper, we describe the theoretical framework to include the blend lines formed under non-LTE conditions, in the radiative transfer equation, and the numerical techniques to solve it. The properties of a blend line having an intrinsic polarization of its own and its interaction with the main line are discussed. The results of our analysis show that the influence of the blend lines on the main spectral lines, though small in the present context, is important and needs to be considered when interpreting the polarized spectral lines in the second solar spectrum.

Key words: line: profiles – polarization – radiative transfer – scattering – Sun: atmosphere.

1 INTRODUCTION

Historically, a survey of the linear polarization arising due to coherent scattering, carried out by Stenflo et al. (1983) over the wavelength range 3165–4230 Å of the solar spectrum, revealed the nature and influence of the blend lines on the second solar spectrum (the polarized solar spectrum that is produced by scattering processes). They introduced an empirical relation between the intensity and polarization profiles of intrinsically unpolarized lines. Based on this model, they could obtain a good determination of both the zero-point of the polarization scale and the level of continuum polarization, and could further use the model to remove the effect of the depolarizing blend lines in the second solar spectrum. The high-resolution recording of the second solar spectrum by Stenflo & Keller (1996, 1997) and the atlas of Gandorfer (2000, 2002, 2005) also explicitly showed the importance of blend lines and the polarizing continuum.

The highly structured second solar spectrum is characterized by a polarized background continuum on which both intrinsically polarizing and depolarizing blend lines are superposed. While a relative polarimetric precision of $10^{-5}$ can routinely be achieved in current imaging Stokes polarimetry, a direct observational determination of the zero-point of the polarization scale is not possible with comparable accuracy. Instead the zero-point, which is needed to convert the observed relative polarizations to absolute polarizations, has to be determined by theoretical considerations based on the expected polarization shapes of the depolarizing blend lines. For this reason, the blend lines are of fundamental importance for all observational and theoretical work with the second solar spectrum. The blend line model that was proposed in Stenflo et al. (1983) was later applied in a somewhat extended way in Stenflo (2005) for the empirical determination of the polarization of the continuous spectrum based on Gandorfer’s atlas.

The theoretical modelling of the line polarization in the second solar spectrum is always associated with incorporating the depolarizing blend lines, as they invariably affect the shapes of the polarized main lines. Blend lines are usually treated by assuming that they are formed in local thermodynamic equilibrium (LTE) conditions, thereby ignoring their own intrinsic polarization. When blend lines are treated in LTE, polarized line and continuum photons are removed due to larger absorption within the line, causing a depolarization of the main line and the continuum (see Fluri & Stenflo 1999, 2001). A theoretical study by Fluri & Stenflo (1999) on the depolarizing blend lines in the visible solar spectrum showed that the relative intensity and relative polarization profiles defined with respect to the continuum are approximately proportional to each other, with a proportionality constant that varies with angle, wavelength and line strength.

The blend lines can be treated in LTE if their height of formation corresponds to collision-dominated layers. If they are formed in low-density layers, it may be necessary to treat them as being formed under non-local thermodynamic equilibrium (NLTE) conditions. Analysis by Fluri & Stenflo (2003) also showed that the
depolarization of continuum by absorbing blend lines rapidly decreases with increasing height of formation, while the depolarization by scattering blend lines increases with height of formation. These calculations were performed with realistic model atmospheres and hypothetical or real line profiles formed in these atmospheres. Although the blend lines were treated in NLTE, their intrinsic polarizability factor was assumed to be zero (namely the blend lines do not have linear polarization of their own). When they have a non-zero polarizability factor, it may also become necessary to treat their intrinsic polarization, to represent their contribution to the polarization profile of the main line (depolarization or repolarization).

In this paper, we consider blend lines with intrinsic polarization, which occur in the wings of the main lines. The number density of strongly polarizing lines is modest in the visible part of the second solar spectrum. Therefore, the conditions that we have envisaged in this paper (namely the proximity of polarizing main and blend lines) are not often realized. However, as we go down in the UV, the second solar spectrum gets increasingly more crowded with strongly polarizing lines. There we can find several good examples of polarizing blend lines [see the UV atlas of Gandorfer (2005)]. Therefore, the theoretical studies presented in this paper become relevant in the analysis of the scattering polarization of the lines in the UV region of the second solar spectrum.

Blend lines belonging to different elements ‘interact’ with the main line of interest through radiative transfer effects (i.e. they couple to the main line through the opacity distribution and multiple scattering). The strength of this interaction depends on their wavelength separation. This interaction is an example of incoherent superposition of the lines. On the other hand, the interactions between the line components of multiplets like Ca II H & K, Na I D1 & D2, Cr I 5204–5208 multiplet, etc. represent coherent superposition of lines. These interactions between the lines arise due to quantum interference between fine structure states or hyperfine structure states of an atom (Stenflo 1997; Smidta et al. 2011).

In Section 2, we describe the formulation of the relevant transfer equation. In Section 3, the numerical methods used to solve the transfer equation are discussed. In Section 4, we present the results of this so-called ‘multiline transfer’ with resonance and Hanle scattering. Concluding remarks are presented in Section 5.

2 FORMULATION OF THE TRANSFER PROBLEM

We consider a 1D, plane-parallel, static, isothermal atmosphere with homogeneous layers. Only one polarizing blend line in the wings of the main line is considered. Line polarization arises from resonance scattering processes (both in the main and the blend lines). In the presence of magnetic fields, this polarization is modified by the Hanle effect. It is sufficient to describe the radiation field by the Stokes vector $I = (I, Q, U)^T$ because we limit our attention to linear polarization, which, for the weak-field Hanle effect, is fully decoupled from circular polarization. In the presence of a magnetic field and the blend line polarization, the total source vector in the Stokes vector basis may be written as

\[
S(\tau, \lambda, \Omega) = \frac{k_1 \phi_1(\lambda) S_1(\tau, \lambda, \Omega) + k_b B(\lambda) U}{k_1 \phi_1(\lambda) + k_b \phi_b(\lambda) + \sigma_{bc} + k_c} + \frac{k_1 \phi_1(\lambda) S_1(\tau, \Omega) + \sigma_{bc} S_0(\tau, \lambda, \Omega)}{k_1 \phi_1(\lambda) + k_b \phi_b(\lambda) + \sigma_{bc} + k_c}.
\]

where $k_1$ and $k_b$ are the frequency-integrated main and blend line absorption coefficients, respectively, $\sigma_{bc}$ and $k_c$ are the continuum scattering and absorption coefficients. $\phi_1$ and $\phi_b$ denote the absorption profiles for the main and the blend lines. Throughout this paper, the symbols ‘1’ and ‘b’ stand for the ‘main’ line and the ‘blend’ line, respectively. $\tau$ is the total optical depth scale defined by

\[
dr = -[k_1 \phi_1(\lambda) + k_b \phi_b(\lambda) + \sigma_{bc} + k_c] \, dr.
\]

The ray direction $\Omega$ is defined by its polar angles ($\theta, \chi$) with respect to the atmospheric normal. In equation (1) $U = (1, 0, 0)^T$. The line source vectors for the main and the blend lines are $S_1$ and $S_0$, respectively. The continuum scattering source vector is denoted by $S_0$. They are given by

\[
S_1(\tau, \lambda, \Omega) = e_1 B(\lambda) U + (1 - e_1) \int \frac{d\Omega}{4\pi} \frac{R(\lambda, \lambda', \Omega', B)}{\phi_1(\lambda)} I(\tau, \lambda', \Omega),
\]

\[
S_0(\tau, \lambda, \Omega) = e_b B(\lambda) U + (1 - e_b) \int \frac{d\Omega'}{4\pi} \frac{R(\lambda, \lambda', \Omega', B)}{\phi_b(\lambda)} I(\tau, \lambda', \Omega)
\]

and

\[
S_0(\tau, \lambda, \Omega) = \oint d\Omega' \frac{d\Omega'}{4\pi} \frac{P(\Omega, \Omega')}{I(\tau, \lambda', \Omega')} \delta(\lambda - \lambda'),
\]

where $e_1$ and $e_b$ are the thermalization parameters for the main and blend lines respectively, and $B(\lambda)$ is the Planck function. For simplicity, $B(\lambda)$ is taken as the same for both the main and the blend lines. The continuum is assumed to be scattering coherently through Rayleigh and Thomson scattering. $P(\Omega, \Omega')$ is the Rayleigh phase matrix (see e.g. Chandrasekhar 1950). The redistribution matrix $R(\lambda, \lambda', \Omega, \Omega'; B)$ is factorized in the form:

\[
R(\lambda, \lambda', \Omega, \Omega'; B) = R(\lambda, \lambda') P(\Omega, \Omega'; B),
\]

where $\Omega' = (\theta', \chi')$ denotes the incoming-ray direction and $B$ is the vector magnetic field. $R(\lambda, \lambda')$ is the angle averaged redistribution function of Hummer (1962). The quantity $P(\Omega, \Omega'; B)$ is the Hanle phase matrix (see Stenflo 1978, Landi Degl’Ignocenti & Landi Degl’Ignocenti 1988). For clarity, we present the equations for a simple version of the redistribution matrix $R$. In particular, we neglect depolarizing elastic collisions and consider only pure type II scattering in the main line. The blend line is assumed to be scattering according to either partial frequency redistribution (PRD) or complete frequency redistribution (CRD). An exact treatment of collisions according to the approximation level III of Bommier (1997b) can easily be incorporated into the present formalism. Calculations using such physically realistic redistribution matrix are presented in Section 4.8.

For isothermal slab models, we introduce the parameters

\[
\frac{\beta_c}{k_c} = \frac{k_1}{k_c}, \quad \frac{\beta_b}{k_c} = \frac{k_0}{k_c}, \quad \frac{\beta_{bc}}{k_c} = \frac{\sigma_{bc}}{k_c}.
\]

Further, we work in the irreducible basis (see Frisch 2007), where the source vector depends only on $\tau$ and $\lambda$. In this basis, using the Hanle phase matrix elements in the atmospheric reference frame, it is easy to show that the total and the line source vectors have the form:

\[
S(\tau, \lambda) = \frac{[\beta_c \phi_1(\lambda) + \beta_b \phi_b(\lambda) + \beta_{bc}] S_1(\tau, \lambda) + B(\lambda) U}{\beta_c \phi_1(\lambda) + \beta_b \phi_b(\lambda) + \beta_{bc} + 1}.
\]
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and

\[
\mathbf{S}_L(\tau, \lambda, \mu) = \frac{\beta_c \phi_c(\lambda) \epsilon_l + \beta_b \phi_b(\lambda) \epsilon_b}{\beta_c \phi_c(\lambda) + \beta_b \phi_b(\lambda) + \beta_{sc}} B(\lambda) \mathbf{U} + \int_{-1}^{+1} \frac{d\mu'}{2} \int_{0}^{\infty} d\lambda' \left[ \frac{\mathbf{N}^R(\mathbf{B}) \beta_b(1 - \epsilon_b) \mathbf{W}^R(\lambda, \lambda')}{\beta_c \phi_c(\lambda) + \beta_b \phi_b(\lambda) + \beta_{sc}} \right. \\
\left. + \frac{\mathbf{N}^R(\mathbf{B}) \beta_b(1 - \epsilon_b) \mathbf{W}^R(\lambda, \lambda')}{\beta_c \phi_c(\lambda) + \beta_b \phi_b(\lambda) + \beta_{sc}} \right] \times \Psi(\mu') \mathbf{I}(\tau, \lambda', \mu'). \tag{7}
\]

Here \( \mathbf{U} = (1, 0, 0, 0, 0, 0)^T \), \( \mathbf{R}^R(\lambda, \lambda') \) is either given by \( R^R_0(\lambda, \lambda') \) or CRD. Note that we have combined the line source vectors for both the main and the blend lines, as well as the continuum scattering source vector, in a single expression. This allows us to apply the approximate lambda iteration (ALI) method of solution based on the frequency by frequency (FFB) technique to compute the line source vector corrections. \( \Psi(\mu') \) is the Rayleigh phase matrix in the irreducible basis, \( \mathbf{N}^R(\mathbf{B}) \) and \( \mathbf{N}^Q(\mathbf{B}) \) are the Hanle phase matrices in the irreducible basis for the main and blend lines, respectively. Expressions for these can be found in Frisch (2007). Hanle phase matrices for the two lines could be different as they can form at different heights in the atmosphere, with different strength and geometry of the magnetic fields. In the absence of magnetic fields, \( \mathbf{N}^R(\mathbf{B}) \) and \( \mathbf{N}^Q(\mathbf{B}) \) matrices reduce to unity matrix \( \mathbf{E} \). The matrices \( \mathbf{W}^R \) and \( \mathbf{W}^Q \) are diagonal, with \( W^R(\lambda, \lambda) = 1 \) and \( W^Q(\lambda, \lambda) = W^Q_\lambda \), where \( k = 1, 2, 3, 4, 5 \). Here \( W^Q_\lambda \) are called polarizability factors. They depend on the angular momentum quantum numbers of the upper and lower levels. For a normal Zeeman triplet transition \( (J = 0 \rightarrow 1 \rightarrow 0) \), this factor is unity.

The 1D line transfer equation for polarized Hanle scattering problem in the irreducible basis is then given by

\[
\frac{\partial S(\tau, \lambda, \mu)}{\partial \tau} = \mathbf{I}(\tau, \lambda, \mu) - \mathbf{S}(\tau, \lambda). \tag{8}
\]

\( \mathbf{I} \) is the formal six-component vector. Our task is to solve this transfer equation to obtain the Stokes profiles \( I, Q, U \) and \( V \). For this purpose, we use the scattering expansion method (SEM) proposed by Frisch et al. (2009).

3 THE NUMERICAL METHOD OF SOLUTION

Since the solution of the transfer equation by polarized ALI (called PALI) method (see Nagendra et al. 1999) with the FFB technique is computationally expensive (see e.g. Sampoorna, Nagendra & Frisch 2008), we opt for SEM, which is a faster approximate method. It is based on Neumann series expansion of the polarized component of the source vector. We extend this method presented in Frisch et al. (2009) to the problem at hand. The source vector component in the irreducible basis can then be written as

\[
S^K_{Q,L}(\tau, \lambda, \mu, \lambda') = \frac{\beta_c \phi_c(\lambda) \epsilon_l + \beta_b \phi_b(\lambda) \epsilon_b}{\beta_c \phi_c(\lambda) + \beta_b \phi_b(\lambda) + \beta_{sc}} B(\lambda) \delta_{K,0} \delta_{Q,0} + \int_{-1}^{+1} \frac{d\mu'}{2} \int_{0}^{\infty} d\lambda' \sum_{Q'} \frac{R^K_{Q,Q'}(\lambda, \lambda')}{\phi(\lambda)} \times \Psi^{K,K}_{Q'}(\mu') I^K_{Q'}(\tau, \lambda', \mu'), \tag{9}
\]

where

\[
\phi(\lambda) = \beta_c \phi_c(\lambda) + \beta_b \phi_b(\lambda) + \beta_{sc} \tag{10}
\]

and

\[
R^K_{Q,Q'}(\tau, \lambda, \lambda') = \frac{\beta_c(1 - \epsilon_b)R^K(\lambda, \lambda') W^K_\lambda}{\phi(\lambda)} \frac{R^0_{Q,Q'}(\lambda, \lambda')}{\phi(\lambda)} + \beta_{sc} \delta(\lambda - \lambda') \frac{\delta_{Q,Q'}}{\phi(\lambda)} \quad \delta_{Q,Q'} \delta_{Q',0}. \tag{11}
\]

In the solar atmosphere, the degree of anisotropy is of the order of few per cent. Thus the degree of linear polarization that arises due to Rayleigh scattering is small. In other words, for the calculation of Stokes \( I \), one can neglect the contribution from the linear polarization \( (Q, U) \) to \( I \) to a good approximation. Therefore, the dominant contribution to Stokes \( I \) comes from the component \( I^0_0 \).

The corresponding source vector component, neglecting the \( K \neq 0 \) terms, is given by

\[
S^0_0 = \frac{\beta_c \phi_c(\lambda) \epsilon_l + \beta_b \phi_b(\lambda) \epsilon_b}{\beta_c \phi_c(\lambda) + \beta_b \phi_b(\lambda) + \beta_{sc}} B(\lambda) + \int_{-1}^{+1} \frac{d\mu'}{2} \int_{0}^{\infty} d\lambda' \frac{R^0_{0,Q}(\tau, \lambda', \mu')}{\phi(\lambda)}. \tag{12}
\]

Here \( S^0_0 \) stands for approximate value of \( S^0_0 \). It represents the solution of a NLTE unpolarized radiative transfer equation. We calculate it using the ALI method of solution with the FFB technique (see Paletou & Auer 1995).

Retaining only the contribution from \( I^0_0 \) on the right-hand side of \( K = 2 \) component of \( S^0_0 \) in equation (9), we obtain the single scattering approximation for each component \( S^2_{Q,L} \) as

\[
[S^2_{Q,L}(\tau, \lambda)]^{(1)} \simeq \int_{-1}^{+1} \frac{d\mu'}{2} \int_{0}^{\infty} d\lambda' \frac{R^2_{Q,Q'}(\lambda, \lambda')}{\phi(\lambda)} \times \Psi^{2,2}_{Q'}(\mu') I^2_{Q'}(\tau, \lambda', \mu'). \tag{13}
\]

The superscript (1) stands for single scattering. The single scattered polarization radiative field \( I^2_{Q'}(\tau, \lambda') \) is calculated using a formal solver. This solution is used as a starting point to calculate the higher order scattering terms. Thus, the iterative sequence at order \( n \) is

\[
[S^2_{Q,L}(\tau, \lambda)]^{(n)} \simeq [S^2_{Q,L}(\tau, \lambda)]^{(1)} + \int_{-1}^{+1} \frac{d\mu'}{2} \int_{0}^{\infty} d\lambda' \times \sum_{Q''} \frac{R^2_{Q,Q''}(\lambda, \lambda')}{\phi(\lambda)} \Psi^{2,2}_{Q''}(\mu') [I^2_{Q''}(\tau, \lambda', \mu')]^{(n-1)}. \tag{14}
\]

The iteration is continued until the maximum relative change in surface polarization becomes less than the convergence criteria of \( 10^{-8} \).

4 RESULTS

When modelling the specific lines of the second solar spectrum, the blend lines are generally treated in LTE. In that case, the blend usually depolarizes the main line polarization. In this section, we consider a polarizing blend line that is formed both in the presence and absence of magnetic fields. We present the dependence of the main line polarization on the blend line polarizability, its separation from the main line and its strength. Effects of variation of \( T \) – the optical thickness of the isothermal slab, \( \epsilon_b \) – the thermalization parameter of the blend line and \( \beta_c \) – the ratio of the background absorbing continuum opacity to the main line opacity, on the main line polarization, are considered. The role played by the Hanle effect and collisions are also discussed. Finally, a brief discussion on the behaviour of the scattering continuum is given.
**Standard model.** In the case of an isothermal atmosphere, the emergent intensity and polarization spectra resemble closely the realistic situation, for the following model parameters. A self-emitting slab of optical thickness $T = 10^6$ is considered. The ratio of background absorbing continuum opacity to the main line opacity $\beta_c = 10^{-2}$, the main line strength $\beta_\lambda = 10^7$, the blend line strength $\beta_b = 5 \times 10^7$, and the continuum scattering coefficient $\beta_sc = 0$. The thermalization parameters are $\epsilon_l = 10^{-2}$ and $\epsilon_c = 5 \times 10^{-2}$. The damping parameters of the main and blend lines are $2 \times 10^{-3}$ and $10^{-4}$, respectively. Both the lines scatter according to pure $K_\parallel$ in the absence of magnetic fields. $W_2$ of both the lines are assumed to be unity. The main line is centred at 5000 Å and the blend line at 5001 Å. The Doppler width is 0.025 Å for both the lines. We refer to this model as the ‘standard model’ and the Stokes profiles for this model are shown as ‘thin solid lines’ in most of the figures. The vertical solid lines represent the ‘wavelength positions’ of the main and blend lines in the figures presented throughout this paper.

### 4.1 Effect of the polarizability factor $W_2^b$ of the blend line

Fig. 1 shows the emergent intensity and polarization profiles. Initially we treat the blend line to be depolarizing ($W_2^b = 0$) and gradually increase the value of $W_2^b$ until it becomes completely polarizing ($W_2^b = 1$). The variation in $W_2^b$ causes very little or no change in the intensity. As expected, PRD triple peaks in $Q/I$ are clearly visible in the case of the main line. The PRD peaks of the blend line are not seen, since the blend line is assumed to be weaker than the main line. If the blend line has zero intrinsic polarization (see the dotted line in the inset), then the wing polarization of the main line is reduced at the core position of the blend line. The extent of depolarization depends on the blend line strength. In the case presented in Fig. 1 (where the blend line is not too strong), we still see a significant depolarization at the blend line core. When the blend line has a non-zero intrinsic polarizability, a peak at the wavelength position of the blend line is observed. As expected, with an increase in $W_2^b$ the polarization of the blend line increases from 0 per cent (when $W_2^b = 0$) to nearly 2 per cent (when $W_2^b = 1$). Since the blend line is very weak, the polarization of the main line is insensitive to the changes in the polarizability factor of the blend line outside the narrow core region of the blend line.

### 4.2 Effect of separation of the blend line from the main line

The relevant results are shown in Fig. 2. The main line is kept fixed at 5000 Å and the position of the blend line is varied. The influence of the blend line on the main line remains limited to the core region and the immediate surroundings of the blend line, as it does not have a significant wing opacity due to its weakness. The ratio of the blend line opacity to the main line opacity increases as the blend line is shifted away from the main line centre (because the main line opacity is relatively small in the far wings). This change in the opacity ratio makes the blend line intensity profile more and more deep, along with corresponding increase in $Q/I$, at the wavelength positions of the blend line. As the line separation increases, the two lines are weakly coupled by transfer effects, eventually becoming completely independent. The profiles computed by treating the blend line in PRD or in CRD are similar because the blend line is assumed to be weak.

### 4.3 Effect of the blend line strength

The strength of the blend line is varied from $\beta_c = 5 \times 10^2$ to $5 \times 10^4$. The blend line is positioned 1 Å away from the main line.
line. The emergent intensity and polarization profiles are shown in Fig. 3. As $\beta_b$ increases, the blend line optical thickness increases resulting in relatively larger heights of the blend line PRD wing peaks in intensity. In the $Q/I$ panel, the PRD peaks of the blend line become more and more prominent as its strength increases. This occurs because of the enhanced scattering opacity as a result of which the near wing polarization of the blend line increases. The Stokes profiles computed treating the blend line in CRD are not significantly different from those computed using PRD as the blend line strength continues to be smaller than the main line strength.

4.4 Effect of variation of optical thickness $T$ of the isothermal slab

The effect of a polarizing blend line on the main line polarization profile with the variation in the total optical thickness $T$ of the isothermal slab is shown in Fig. 4. The blend line is much weaker than the main line and scatters according to pure $R_b$. As $T$ increases, the main line changes from a self-reversed emission line to an absorption line (see the intensity panel in Fig. 4). As the main line core is already saturated, the effect of increase in $T$ is felt only in the line wings. As for the blend line, when $T = 10^3$, a weak line is formed because of the smaller number of main line photons available for scattering. As $T$ increases, the blend line starts to show up prominently in intensity.

The main line polarization profile shows a typical triple peak structure (due to PRD mechanism) when $T = 10^3$. However, the main line near wing PRD peak changes over from negative maxima to positive maxima, as $T$ increases. This has a direct correlation with the behaviour of the Stokes $I$ profile in the region of near wing maxima. The change in sign is indicative of a switchover from limb brightening to the limb darkening of the radiation field at heights where the monochromatic optical depths corresponding to the near wing maxima are unity. When $T = 10^5$ the blend line shows a double-peak structure in $Q/I$, although it is weaker in intensity. The polarization is quite strong, as the blend line is assumed to be polarizing with $W^b = 1$. As $T$ increases, the double-peak structure changes over to a single-peak structure. Away from the blend line centre, the $Q/I$ profiles of the blend line smoothly merge with the main line polarization profiles.

4.5 Effect of variation of the destruction probability $\epsilon_b$

For the ‘standard model’ considered in this paper, the variation in $\epsilon_b$ does not produce significant changes in the $(I, Q/I)$ profiles of the main line. With a decrease in the value of $\epsilon_b$, the blend line depth increases (going from LTE-like to NLTE-like expected behaviour), which saturates for $\epsilon_b < 10^{-3}$ (figure not shown). In $Q/I$, the blend line core peak increases with the decrease in $\epsilon_b$, as the blend line becomes more and more scattering dominated. Further, for $\epsilon_b < 10^{-3}$, the blend line core peak in $Q/I$ saturates, an effect discussed in Faurobert (1988) for the single line case.

4.6 Effect of variation of the continuum absorption parameter $\beta_c$

The ratio of the background continuum absorption opacity to the main line opacity, $\beta_c$, is varied from $10^{-3}$ to $10^{-9}$ (see Fig. 5). The main line strength ($1/\beta_c$) correspondingly changes. This variation of $\beta_c$ influences the Stokes profiles of both the lines. The intensity profiles become narrow and shallow with the increase in the value of $\beta_c$. This is because the continuum progressively influences the inner parts of the main line profile as $\beta_c$ increases. The main line...
which was a pure absorption line when $\tilde{\beta}_c = 10^{-3}$ (the dotted line) becomes a self-reversed emission line when $\tilde{\beta}_c = 10^{-9}$ (dotted-dashed line). The decrease in wing intensity is due to a decrease in the continuum optical thickness ($T^c = \tilde{\beta}_c T$) as $\tilde{\beta}_c$ varies from $10^{-3}$ to $10^{-9}$. The blend line intensity profile also changes from a strong absorption line to a weak emission line as $\tilde{\beta}_c$ changes from $10^{-3}$ to $10^{-9}$.

The main line $Q/I$ core amplitude is not very sensitive to $\tilde{\beta}_c$ unless $\tilde{\beta}_c$ is sufficiently large (see the dotted line). The main line near wing PRD peak, as well as far wing polarization, decreases in magnitude as $\tilde{\beta}_c$ increases, due to the predominance of the unpolarized continuum. The $Q/I$ profiles at the core of the blend line nearly coincide for $\tilde{\beta}_c = 10^{-3}, 10^{-5}$ and $10^{-7}$. However, when $\tilde{\beta}_c = 10^{-9}$, the blend line acts like a depolarizing line in the wing of the main line, in spite of $W_2^b$ being unity. This is possibly because the blend line in this case is a weak emission line, whose polarization profiles are characteristically different from those of absorption lines.

4.7 Multiline transfer in the presence of magnetic fields

The vector magnetic field $B$ is parametrized through $(\gamma_B, \theta_B, \phi_B)$, with $\gamma_B = g_i \alpha_i / \Gamma_K$, where $g_i$ is the upper level Lande-$g$ factor, $\alpha_i$ the Larmor precession rate and $\Gamma_K$ the damping rate (inverse lifetime) of the excited state (see e.g. Stenflo 1994). The magnetic field orientation represented by $\theta_B$ and $\phi_B$ are defined with respect to the atmospheric normal. The $\gamma_B$ for the main line is fixed as unity. The $\gamma_B$ of the blend line is varied from 0 to 10. Fig. 6 shows the profiles for the two-line system in the presence of magnetic fields. The blend line shows similar effects on the main line both in the presence and absence of magnetic fields, for the model parameters used in this section. The main and blend line intensities are unaffected. The magnitude of $Q/I$ in the central peak of the blend line reduces with an increase in the value of its $\gamma_B$. This is the typical effect of magnetic fields, namely the Hanle effect which is operative in the core regions of the two lines. Stokes $U$ which was zero for Rayleigh case is generated by the Hanle effect and hence characteristic core peaks are seen in the $U/I$ panel. The depolarization in the core region of the blend line due to Hanle effect causes a corresponding increase in $U/I$.

4.8 Effect of elastic collisions

It is well known that the Hanle effect operates efficiently in the line core (within a few Doppler widths) and disappears in the line wings (Omont, Smith & Cooper 1973). The functional form of this frequency dependence of the Hanle effect is presented in Stenflo (1998). To account for this frequency dependence of the Hanle effect in numerical computations, we introduce the 1D cut-off approximation. Fig. 6 presented in Section 4.7 is computed using the 1D cut-off approximation, which implies the use of Hanle phase matrices. The exact collisional PRD theory of Hanle effect as well as the approximations leading to these 2D domain based PRD formulation are developed by Bommier (1997a,b).
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The effect of depolarizing elastic collisions. The magnetic field parameters are \((\gamma_B, \theta_B, \phi_B) = (1.30, 0^\circ, 0^\circ)\) for both the lines. Other model parameters are the same as in the ‘standard model’.

The strength of elastic collisions is specified through \(\Gamma_E/\Gamma_R\), where \(\Gamma_E\) denotes the elastic collisional rate and \(\Gamma_R\) the radiative de-excitation rate. The values of \(\Gamma_E/\Gamma_R\) chosen by us cover the situations ranging from the absence of elastic collisions (pure \(\Gamma_R\)) to the presence of strong elastic collisions. Depolarizing collision rates are given by \(D = c(\Gamma_E/\Gamma_R)\) with \(c = 0.5\) (see Stenflo 1994).

The emergent Stokes profiles are shown in Fig. 7. They refer to the cases where \(\Gamma_E/\Gamma_R\) of both the lines are taken as equal and varied in the same fashion. In all these cases, we see that the elastic collisions do not modify the intensities in the cores of the two lines. This is because in the line core \(\Gamma_R\) behaves more like CRD. In the wings of the two lines, the PRD-like intensity profiles gradually approach the CRD-like behaviour (true absorption line), with an increase in the elastic collision rate \(\Gamma_E/\Gamma_R\). For large values of \(\Gamma_E/\Gamma_R\) (\(=100\)), the line polarization approaches zero (double-dot–dashed line) throughout the line profiles. \(U/I\) profiles also show similar behaviour as the \(Q/I\) profiles in both the lines.

5 CONCLUSIONS

In this paper, we present our detailed studies on the effects of a blend line (polarizing or depolarizing) present in the wings of a main line. Our particular interest is the linear polarization profiles of the main line. We show how theoretically the total source function can be generalized to include a blend line. The same formalism can be extended to deal with the cases where there is more than one blend line. We formulate the radiative transfer equation in the irreducible tensorial basis. This transfer equation is then solved by computing the scalar intensity with the standard FBF iterative technique, and the polarization by a faster, promising method called SEM. In the SEM, the polarized component of the source vector is expanded in Neumann series. Single scattered solution is computed at first, and this solution is then used for calculating the higher order scattering terms. The dependence of the blend line intensity and polarization effects on various parameters like polarizability factor, distance from the main line core, blend line strength, isothermal slab optical thickness, continuum opacity and polarization, magnetic field, and elastic collision rate has been explored in detail.

4.9 Continuum polarization

Fig. 8 shows the effect of continuum polarization on the blend and the main line polarization. The continuum polarization arises due to Thomson scattering on electrons and Rayleigh scattering on atoms and molecules. It is included here through the parameter \(\beta_{sc}\).
The blend lines in the linearly polarized spectrum of the Sun invariably affect the main spectral lines. A knowledge of the way in which this interaction takes place plays an important role in the interpretation of the second solar spectrum. The insight that we have gained through our theoretical studies using isothermal slab models is a first step towards realistic modelling of the second solar spectrum. Such calculations become necessary in a fine analysis of the solar spectrum, and help in our studies of the solar atmosphere. The papers by Anusha et al. (2010, 2011) and Smitha et al. (2012) form the basis for the inclusion of intrinsically polarized blend lines in modelling the second solar spectrum.

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REFERENCES


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