Investigation of the Transient Nature of Thunderstorm Winds from Europe, the United States, and Australia Using a New Method for Detection of Changepoints in Wind Speed Records

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ABSTRACT

This paper investigates the transient characteristics in 41 velocity records of 19 thunderstorm events from around the world—9 from Europe, 9 from the United States, and 1 from Australia. The transient features of thunderstorm winds were examined by introducing an objective method for the detection of changepoints in the time series. The methodology divides velocity records into segments characterized by different statistical properties. The segmentation is based on the following properties of the isolated segments: mean (M) and the standard deviation (SD). This study demonstrated that the maximum velocity during the thunderstorm peak in the events from Europe is typically 2–4 times larger than the mean wind speed before the thunderstorm. The duration of the thunderstorm velocity peak was 2–5 min in approximately 60% of the analyzed records using the M statistic and 5–10 min when analyzed using the SD statistic. Therefore, the velocity fluctuations caused by thunderstorm winds last longer than the abrupt changes in the mean wind field. Similarly, the ramp-up time was longer when the records were analyzed using the SD statistic. The segmentation methodology was tested for different duration of velocity records and using data with different sampling frequencies. The performances of the introduced method were compared against the results of two other segmentation procedures proposed in the literature. One of the practical applications of this method is the physical separation between the thunderstorm and nonthunderstorm components of a wind event.

SIGNIFICANCE STATEMENT

Thunderstorm outflow winds are short-lived phenomena produced by cold downdrafts that originate in thunderstorm clouds. This study analyzes the transient nature of thunderstorm winds from Europe, the United States, and Australia using a segmentation method applied to anemometer velocity records. This segmentation method identifies abrupt changes of mean wind speed and wind fluctuations in the velocity data. This research provides the means of isolating different segments within the thunderstorm wind records in an objective way that is based on rigorous mathematical principles. The proposed method can automatically distinguish thunderstorm from nonthunderstorm winds. The peak velocities in thunderstorm outflows are 2–4 times the mean wind speed before the thunderstorm. The most intense episodes of thunderstorm winds usually last 2–5 min.

1. Introduction

Thunderstorms are severe weather phenomena that produce strong winds close to the surface, as well as potentially large amounts of precipitation, hail, and lightning. Analyzing catastrophic events caused by thunderstorms from 1948 to 1998 across the United States, Changnon (2001) noticed that, while the major damage caused by some thunderstorms was due to either hail, lightning, tornadoes or heavy rain, all of the investigated events were also characterized by strong surface winds. The downdraft-associated winds from thunderstorms are
usually classified into downbursts, gust fronts, rear-flank downdrafts (RFD), and forward-flank downdrafts (FFD). Conceptual models of these downdrafts, their origin, and dynamics can be found in Goff (1976), Lemon and Doswell (1979), Wakimoto (1982), Fujita (1985), Markowski (2002), and Shabbott and Markowski (2006), among others.

The transient nature of thunderstorm winds is best depicted by comparing their velocity records against a steady atmospheric boundary layer (ABL) wind (Fig. 1). While a time-independent mean is sufficient at describing the mean ABL flow (Fig. 1a), the same approach does not apply in the example of thunderstorm winds (Fig. 1b). That is, the classical Reynolds decomposition in which the instantaneous wind speed is decomposed into the time-independent mean and time-dependent, but stationary, residual fluctuations does not hold for thunderstorm winds (Fig. 1b).

Fujita (1985) proposed an algorithm for microburst detection based on the 1-min maximum wind speed $U_{\text{1min}}$, as well as the prepeak mean wind speed $\bar{U}_-$ and postpeak mean wind speed $\bar{U}_+$. The last two ($\bar{U}_-, \bar{U}_+$) are calculated from seven 1-min segments with the first minute before and after the peak excluded from the analysis. The following conditions have to be satisfied simultaneously in order for an event to be classified as microburst: 1) $U_{\text{1min}} > 10 \text{ m s}^{-1}$, 2) $U_{\text{1min}} > \bar{U}_+ + 5$ and $U_{\text{1min}} > \bar{U}_+ + 5$, 3) $U_{\text{1min}} > 1.25 \bar{U}_+$, and $U_{\text{1min}} > 1.25 \bar{U}_-$, and last 4) $\bar{U}_+ \leq 1.5 \bar{U}_-$. Microbursts are downbursts with the horizontal scale of the outflow below 4 km. Fujita (1985) applied this algorithm to over 5 million 1-min wind records and extracted 579 microbursts from Northern Illinois Meteorological Research on Downburst (NIMROD) and Joint Airport Weather Studies (JAWS) field campaigns combined.

Lombardo et al. (2009) put forward a methodology that separates thunderstorm from nonthunderstorm winds in the Automated Surface Observing System (ASOS) hourly data. The ASOS network of automated weather stations covers most of the continental United States (NCEI 1998; Romanic et al. 2018; Lombardo and Ziekar 2019). Due to the coarse time resolution of ASOS data, only the hourly peak values of thunderstorm winds and their duration were obtained and analyzed. Since the method was applied to a large set of data, the research was focused on the underlying statistical distribution of extremes, and the specific time series were not investigated in detail. De Gaetano et al. (2014) implemented an automated procedure for the extraction of three classes of winds from anemometer records: 1) stationary and Gaussian, 2) nonstationary and non-Gaussian, and 3) stationary and non-Gaussian winds. The separation methodology of De Gaetano et al. (2014) was implemented through a sequential algorithm that uses 10-min and 1-h values of various statistical parameters averaged over 1-s and 1-min intervals, respectively. The parameters considered over a 10-min interval were peak and mean velocities ($U_{\text{10min}}$ and $\bar{U}_{\text{10min}}$, respectively), mean wind direction ($\alpha_{\text{10min}}$), gust factor ($G_{\text{10min}} = U_{\text{10min}}/\bar{U}_{\text{10min}}$), turbulence intensity ($I_{\text{10min}}$), skewness ($\gamma_{\text{10min}}$), and kurtosis ($k_{\text{10min}}$), as well as the same set of parameters for a 1-h interval (except for the peak velocity). If $U_{\text{10min}} \geq 15 \text{ m s}^{-1}$, the wind record is a candidate to be a thunderstorm wind. Also, if $G_{\text{10min}} > 1.25$ ($G_{\text{10min ref}}$ is the reference gust factor obtained from ESDU (1993)), then the event is either nonstationary and non-Gaussian (i.e., thunderstorm wind) or stationary and non-Gaussian (i.e., intermediate wind). Last, whether such an event is a thunderstorm or intermediate wind is later determined qualitatively. The methodology proposed in this paper could be used to carry out this separation procedure automatically.

Lombardo (2009), Lombardo et al. (2014), and Gunter et al. (2017) investigated individual time series of thunderstorm winds by analyzing abrupt changes in wind speed and direction records. The segmentation method proposed in Lombardo et al. (2014) identifies abrupt changes in wind speed records by inspecting statistically significant changes in variance using the reverse arrangement test (Kendall et al. 1983) and run test (Brownlee 1965;
In this approach, the beginning of time series is split into eighteen 3-s segments that are subjected to stationarity testing. If the test statistics of the portion of the time series were stationary, the algorithm includes the following 3-s interval and repeats the stationarity testing. This addition of the following (3-s) portion to the previous length of the segment is carried out until one or both tests indicate nonstationarity, which is then deemed as a changepoint in the time series. More recently, Gunter et al. (2017) separated the time series of thunderstorm winds into five segments identified in a 40-s moving average of the wind direction record. Their method is based on identifying the zero-derivative points on either side of the abrupt change in the moving average. The physical justification of the methodology follows from the studies of Goff (1976) and Wakimoto (1982) that describe the kinematics of thunderstorm outflows through a sequence of gust surges each characterized by changes in wind speed and direction (Lompar et al. 2018). In this study, we propose a new segmentation method for the following reasons. First, the reverse arrangement method detects the mean trends only (Pan and Patton 2017), whereas thunderstorm winds are accompanied by changes in wind fluctuations in addition to the changes of mean flow. Second, while the run test inspects a combination of mean and periodic variations in time series, it is not capable of distinguishing between the variations in mean from the variations in periodic trends. The wind direction approach, on the other hand, assumes that thunderstorm velocity records are always associated with changes in wind direction. While this assumption is usually accurate, some anemometer records of thunderstorm winds do not exhibit sudden shifts in wind direction. This point will be demonstrated later in section 3.

This article investigates abrupt changes in wind speed records of thunderstorm winds with the goal of describing the transient nature of the phenomena using objective procedures that facilitates automated data analysis. Since the analysis is currently restricted to wind speed data, the paper aims to identify regions with different wind characteristics by disregarding other atmospheric variables (e.g., temperature, pressure, relative humidity). The introduced methodology, however, could be applied to the analysis of other thunderstorm variables (Lompar et al. 2018), and not just wind speed. The presented method for the analysis of abrupt changes in wind speed records is adopted from signal processing carried out in the fields of speech recognition and brain research (Basseville and Nikiforov 1993; Darkhovski 1994; Lavielle 2005), as well as geosciences (Killick et al. 2012). Recently, Romanic et al. (2019) used the same method to investigate the transient features of experimentally produced downburst-like impinging jets in cross-flow. The method was applied to detect the onset of impinging jets and to find the velocity slowdown in crossflows. The present paper extends the applicability of this methodology to real thunderstorm winds. The current study also investigates a large set of thunderstorm events from different geographic and climate regions around the world. As discussed in Markowski and Richardson (2010) and Lompar et al. (2017), the natural variability in cloud dynamics and microphysics between different climate regions can affect precipitation characteristics and gust front dynamics. For example, most of the U.S. thunderstorm records analyzed in this study are associated with supercell or bow echo thunderstorms, while the majority of European data came from less organized thunderstorms.

The objectively analyzed and segmented thunderstorm velocity records can be of importance in wind energy and wind engineering (Solari 2016). For example, an analysis of structural responses to transient winds is typically carried out via the thunderstorm response spectrum technique (Solari 2016), whereas the structural responses to ABL winds are analyzed through the gust response factor technique (e.g., Davenport 1961; Solari 1983). For the former method, the transient characteristics of the velocity record are expressed through the slowly varying mean and the residual fluctuations (e.g., Choi and Hidayat 2002; Chen and Letchford 2004, 2006; Holmes et al. 2008; Solari et al. 2015; Burlando et al. 2017; Junayed et al. 2019). The proposed methodology in this study can be used to objectively separate different segments in thunderstorm velocity records and therefore facilitate the proper interpretation of obtained structural responses. In addition to the analysis of abrupt changes in the instantaneous wind speed records, this paper also investigates the transient features of the decomposed velocity records. Moreover, an attempt will be made to relate the transient nature of the instantaneous wind speed to the transient features of the decomposed time series (e.g., moving-mean and turbulent fluctuations). These findings will be discussed in terms of surface roughness and anemometer height above ground.

2. Data and methods

2a. Data

Thunderstorm events analyzed in this study come from a variety of sources and cover three continents: Europe (9 events), North America (9), and Australia (1) (Fig. 2). Normalized time series of velocity records from all thunderstorm events are depicted in Fig. 3. Eight events from Europe are from two European Union projects “Wind and Ports” (Solari et al. 2012) and
The goal of these projects was to provide high-quality field measurements of thunderstorm winds in the Ligurian Sea and the north Tyrrhenian Sea, in the Mediterranean (Fig. 2a). Out of eight events from this region (Italy), one event was recorded in Genoa, four were measured in La Spezia, and three were from Livorno. Some events were captured with several anemometers making the number of analyzed velocity records larger than the number of considered events. Anemometer characteristics and site coordinates are provided in Table 1 and Fig. 4. The remaining event from Europe was observed in Finland (Järvi et al. 2007) in 2004 (Table 1 and Figs. 2a, 3a). The events from the Mediterranean correspond to thunderstorm winds with mostly marine exposure, while the thunderstorm in Finland is an event that took place above a forest area. Thus, there is a significant difference between surface roughness of the two regions in addition to the profoundly different climatological characteristics of Finland (northeast Europe) and the Mediterranean coast of Italy (south Europe).

Figure 2b shows the location of thunderstorm winds investigated in the United States. Two events were recorded in Texas—one in Pep (Gunter et al. 2017) (Fig. 3k) and another in Lubbock (Orwig and Schroeder 2007; Holmes et al. 2008) (Fig. 3m). The Lubbock downburst was recorded by twelve anemometers (six masts), which makes it the best-documented downburst event in this study. Another event from the United States was recorded in Syracuse, Kansas, (Gunter and Schroeder 2015) (Fig. 3i) and one at Andrews Air Force Base (AAFB) in Maryland near Washington, D.C. (Fujita 1985) (Fig. 3j). The downburst in Washington, reported by Fujita (1985), is still one of the strongest, if not the strongest, with the peak velocity of 68 m s\(^{-1}\) measured at only 4.9 m above ground. The original anemometer measurements for the Washington event were not available to the authors, and therefore the published time series in Fujita (1985) was digitized (Rohatgi 2019) in order to obtain wind speed data. The uncertainty of digitized data is larger than the real measurements. Another set of five thunderstorm events from the Lubbock–Reese (Texas) was obtained from Lombardo et al. (2014) (Table 1).

The single downburst event from Australia was observed by Sherman (1987) in the suburban region of Brisbane (Figs. 2c, 3s). Similar to Fujita (1985), raw data for this event were not available, and thus the analyzed time series were digitized (Rohatgi 2019). Overall, the number of thunderstorm events investigated in this study is 19, while the total number of velocity records is 42.

Surface roughness characteristics surrounding the measuring site were acquired using Google Earth Pro aerial imagery (Fig. 4). The white dashed circles in Fig. 4 have a diameter of 10 km, with the measurement site being close to the center of the circles. The values of roughness length \(z_0\) were either obtained from the literature that originally published the events (e.g., Lombardo et al. 2014; Gunter et al. 2017) or they were estimated by combining the aerial imagery with the exposure categories provided in ESDU documentation (Table 2). The uncertainties of \(z_0\) estimates using different methods are discussed in Romanic et al. (2018). The anemometers in Italy (GE, LS, and LI; see Table 1 for abbreviations) were located along the coast and the \(z_0\) values are provided for land and sea separately. Other anemometers that recorded the data for this study were installed on land and mostly with uniform roughness surrounding the site (HY, SY, PE, and BR). In the cases of nonuniform roughness, the values of \(z_0\) for different wind direction sectors are provided in Table 2.

b. Changepoints in wind records—Theoretical background

Informally, changepoints represent the points in the time series at which statistical properties of a segment of
observations change. Formally, a changepoint can be defined as follows. Consider a time series $U_t = (U_1, U_2, \ldots, U_T)$ of velocity data (e.g., Fig. 1), where $T$ is the length of the record. That is, $T = T f_s$, where $T$ is the length of time between measurements (s) and $f_s$ is the sampling frequency (Hz). A changepoint occurs if at any time $t \in \{1, 2, \ldots, T - 1\}$ some statistical properties of $\{U_1, \ldots, U_t\}$ and $\{U_{t+1}, \ldots, U_T\}$ change. The above definition can readily be extended to a set of changepoints $m$, with their respective positions being $1 < \tau_1 < \tau_2 < \ldots < \tau_m < T$.

In changepoint detection problems, the goal is to minimize the expression (Lavielle 1999, 2005; Killick et al. 2012):

$$
\sum_{i=1}^{m+1} \left[ \epsilon \left( U_{(\tau_i-1)+1}, U_{\tau_i} \right) + \gamma \right],
$$

where $\epsilon(\cdot)$ is a cost function for segments $\tau_{i-1}$ and $\tau_i$, and $\gamma$ is a penalty function against overfitting. In this study, the minimization is performed using a recursive approach (Killick et al. 2012):
\[ F(s) = \min_{t < s} \left[ F(t) + \mathcal{Z}(U_{(t+1)s}) + \gamma \right], \quad (2) \]

where \( F(s) \) is the optimal partitioning of the points \( U_1, U_2, \ldots, U_s \) \( (s = 1, 2, \ldots, T) \) and \( F(t) \) is the optimal partitioning of the points \( U_1, U_2, \ldots, U_t \), where \( t < s \) and \( F(0) = -\gamma \). Details of this computational algorithm—the so-called optimal partitioning method—are described in Jackson et al. (2005) and Lavielle (2005), with improvements in computational efficiency using the prude exact linear-time method proposed by Killick et al. (2012). Equation (2) shows that the penalty constant \( \gamma \) is a "balance" between the cost function and the number of changepoints. If \( \gamma = 0 \), there is no penalty for adding a changepoint to the model and, therefore, the data are segmented in as many segments as possible (i.e., \( T - 1 \) segments). The relationship between \( \gamma \) and \( m \) depends on the properties of data, and it will be discussed more in section 3d based on the results obtained from the analyzed velocity records. In general terms, however, Eq. (2) demonstrates that a larger \( \gamma \) results in fewer changepoints and thus smaller variance of the model.

Since the variance describes the sensitivity of the model to input data, the small variance also results in a

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**TABLE 1. Thunderstorm wind events investigated in this study.**

<table>
<thead>
<tr>
<th>No.</th>
<th>Country</th>
<th>Anemometer location</th>
<th>Event date</th>
<th>Anemometer height AGL (m)</th>
<th>Sampling frequency (Hz)</th>
<th>Reference study</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>Finland</td>
<td>Hyytiälä (HY)</td>
<td>3 Jul 2004</td>
<td>73</td>
<td>10.4</td>
<td>Järvi et al. (2007)</td>
</tr>
<tr>
<td>2</td>
<td>Italy</td>
<td>Genoa (GE)</td>
<td>30 Sep 2012</td>
<td>61.4</td>
<td>10</td>
<td>De Gaetano et al. (2014); Burlando et al. (2017, 2018)</td>
</tr>
<tr>
<td>3</td>
<td>Italy</td>
<td>La Spezia (LS)</td>
<td>5 Jun 2011</td>
<td>15.5</td>
<td></td>
<td></td>
</tr>
<tr>
<td>4</td>
<td></td>
<td></td>
<td>11 Apr 2012</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>5</td>
<td></td>
<td></td>
<td>19 Apr 2012</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>6</td>
<td></td>
<td></td>
<td>25 Oct 2011</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>7</td>
<td>Livorno (LI)</td>
<td></td>
<td>1 Oct 2012</td>
<td>20</td>
<td></td>
<td></td>
</tr>
<tr>
<td>8</td>
<td></td>
<td></td>
<td>4 Sep 2011</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>9</td>
<td></td>
<td></td>
<td>11 Jun 2011</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>10</td>
<td></td>
<td></td>
<td>26 Oct 2012</td>
<td>75</td>
<td></td>
<td></td>
</tr>
<tr>
<td>11</td>
<td></td>
<td></td>
<td>26 Oct 2012</td>
<td>20</td>
<td></td>
<td></td>
</tr>
<tr>
<td>12</td>
<td></td>
<td></td>
<td>11 Jun 2011</td>
<td>2.25</td>
<td>5</td>
<td>Gunter and Schroeder (2015); Gunter et al. (2017)</td>
</tr>
<tr>
<td>13</td>
<td>United States</td>
<td>Syracuse (SY)</td>
<td>24 Aug 2007</td>
<td>10</td>
<td>30</td>
<td>Lombardo et al. (2014)</td>
</tr>
<tr>
<td>14</td>
<td></td>
<td></td>
<td>21 May 2008</td>
<td>10</td>
<td>30</td>
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<tr>
<td>15</td>
<td></td>
<td></td>
<td>19 Jun 2008</td>
<td>10</td>
<td>30</td>
<td></td>
</tr>
<tr>
<td>16</td>
<td></td>
<td></td>
<td>4 Jun 2009</td>
<td>17</td>
<td>30</td>
<td></td>
</tr>
<tr>
<td>17</td>
<td></td>
<td></td>
<td>14 May 2008</td>
<td>2.5</td>
<td>5</td>
<td></td>
</tr>
<tr>
<td>18</td>
<td></td>
<td></td>
<td>1 Aug 1983</td>
<td>4.9</td>
<td>0.1</td>
<td>Fujita (1985)</td>
</tr>
<tr>
<td>19</td>
<td>Australia</td>
<td>Sydney (SY)</td>
<td>5 Nov 1977</td>
<td>58</td>
<td>1.025</td>
<td>Sherman (1987)</td>
</tr>
<tr>
<td>20</td>
<td></td>
<td></td>
<td>104</td>
<td>1.14</td>
<td></td>
<td></td>
</tr>
<tr>
<td>21</td>
<td></td>
<td></td>
<td>10</td>
<td>0.196</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
more substantial bias of the model. So, in other words, 
$g$ is a trade-off between bias and variance of $F(s)$.

The next step is to describe the cost functions used in this study. Let us again consider the time series $U_t = (U_1, U_2, \ldots, U_T)$ and, supported by the central limit theorem, let us further assume that the time series is a result of an independent normal stochastic process (Chen and Gupta 2012). Therefore, $U_t$ can be modeled as

$$U_t = c_{0_j} + c_{1_j}t_i + e_i,$$

(3)

where $t_i (i = 1, 2, \ldots, T)$ is the time, $e_i$ is the noise (error term with zero mean and unit variance), and $c_{0_j}$ and $c_{1,j}$ are the linear regression coefficients of the segment $j$. In other words, $U_t$ is a subset of a random variable $U_i$ whose Gaussian probability density function $f_{U_i}$ is

$$f_{U_i}(U_i|t_i, c_{0_j}, c_{1_j}, \sigma_j^2) = \frac{1}{\sqrt{2\pi\sigma_j^2}} \exp \left[ -\frac{(U_i - c_{0_j} - c_{1_j}t_i)^2}{2\sigma_j^2} \right],$$

(4)
where $\sigma^2$ is the variance. Invoking the assumption of data independency once more, the likelihood function $L$ of all data $U_i = (U_1, U_2, \ldots, U_T)$ is the product of the $f_{i|j}$:

$$
L(C_0, C_1, \sigma^2) = \prod_{j=1}^{m+1} \prod_{i=\tau_{j-1}+1}^{\tau_j} \frac{1}{\sqrt{2\pi\sigma_j^2}} \exp \left[ -\frac{(U_j - c_{0j} - c_{1j}\bar{t}_j)^2}{2\sigma_j^2} \right], \quad (5)
$$

where $C_0 = (c_{01}, c_{02}, \ldots, c_{0,m+1})$, $C_1 = (c_{11}, c_{12}, \ldots, c_{1,m+1})$, and $\sigma^2 = (\sigma_{01}^2, \sigma_{02}^2, \ldots, \sigma_{m+1}^2)$. The unbiased maximum likelihood estimates of the unknown parameters in Eq. (5) are the solution to the following system of algebraic equations (Neter et al. 1996):

$$
\hat{C}_{0j} = \bar{U}_j - \hat{C}_{1j}\bar{t}_j, \quad (6)
$$

$$
\hat{C}_{1j} = \frac{\sum_{i=\tau_{j-1}+1}^{\tau_j} (t_i - \bar{t}_j)^2(U_i - \bar{U}_j)^2}{\sum_{i=\tau_{j-1}+1}^{\tau_j} (t_i - \bar{t}_j)^2}, \quad (7)
$$

$$
\hat{\sigma}_j^2 = \frac{1}{\tau_j - \tau_{j-1} - 2\sum_{i=\tau_{j-1}+1}^{\tau_j} (U_i - c_{0j} - c_{1j}\bar{t}_j)^2}, \quad (8)
$$

where $\bar{t}_j = 1/\tau_j - \tau_{j-1} - \sum_{i=\tau_{j-1}+1}^{\tau_j} t_i$ and $\bar{U}_j = 1/\tau_j - \tau_{j-1} \sum_{i=\tau_{j-1}+1}^{\tau_j} U_i$. It is beneficial to introduce twice the negative log-likelihood function as a choice of the cost function in order to directly add segment costs to the overall cost (Horvath 1993). Therefore, we cast Eq. (5) in the form:

$$
\mathcal{L} = -2\ln[L(\hat{C}_{0j}, \hat{C}_{1j}, \hat{\sigma}_j^2)], \quad (9)
$$

or after expanding the right-hand side:

$$
\mathcal{L} = (U_{(\tau_j-1)} - \tau_{j-1}) \ln(2\pi\sigma_j^2) + 2 \sum_{i=\tau_{j-1}+1}^{\tau_j} \frac{(U_i - \hat{c}_{0j} - \hat{c}_{1j}\bar{t}_j)^2}{2\sigma_j^2}. \quad (10)
$$

Equation (10) represents the likelihood linear regression cost function. By setting $\hat{c}_1 = 0$, we retrieve the likelihood mean cost function:

$$
\mathcal{L} = -\ln(2\pi\sigma_j^2) + 2 \sum_{i=\tau_{j-1}+1}^{\tau_j} \frac{(U_i - \hat{c}_{0j})^2}{2\sigma_j^2}. \quad (11)
$$

When analyzing the abrupt changes in the mean wind speed, the standard deviation $\sigma$ (square root of the variance) is a fixed constant in the above expression. Similarly, when investigating the changepoints through $\sigma$ the mean is a fixed constant (Chen and Gupta 2012). Knowing that constant terms do not contribute to the result of minimization or maximization problems, the exact value of any constant is irrelevant. Changepoints in the mean and standard deviation statistics are hereafter referred to as the M and SD changepoints, respectively.

c. Velocity decomposition

In contrast to the classical Reynolds decomposition that applies to stationary ABL winds (Fig. 1), transient wind records are usually analyzed via (e.g., Choi and Hidayat 2002; Chen and Letchford 2004; Holmes et al. 2008; Solari et al. 2015; Burlando et al. 2017):

$$
U(t) = \bar{U}(t) + U'(t), \quad (12)
$$

where $\bar{U}(t)$ is the slowly varying wind speed associated with the low-frequency content of $U(t)$, and $U'(t)$ is the residual fluctuations. By further decomposing the residual fluctuations as

<table>
<thead>
<tr>
<th>Site location</th>
<th>Land</th>
<th>Water</th>
</tr>
</thead>
<tbody>
<tr>
<td>Genoa, Italy</td>
<td>4 (rugged hilly terrain)</td>
<td>10^{-3} (large expansion of water)</td>
</tr>
<tr>
<td>La Spezia, Italy</td>
<td>4 (rugged hilly terrain)</td>
<td>10^{-3} (large expansion of water)</td>
</tr>
<tr>
<td>Livorno, Italy</td>
<td>0.05 (farmland) to 0.6 (center of small towns)</td>
<td>10^{-3} (large expansion of water)</td>
</tr>
<tr>
<td>Hyytiälä, Finland</td>
<td>0.5 (fairly level wooded country)</td>
<td>—</td>
</tr>
<tr>
<td>Syracuse, KS</td>
<td>0.01 (fairly level grass plains)</td>
<td>—</td>
</tr>
<tr>
<td>Pep, TX</td>
<td>0.01 (fairly level grass plains) (Gunter et al. 2017)</td>
<td>—</td>
</tr>
<tr>
<td>Lubbock, TX</td>
<td>0.01 (fairly level grass plains) (Lombardo et al. 2014)</td>
<td>—</td>
</tr>
<tr>
<td>Washington, D.C.</td>
<td>4 × 10^{-3} (airport) to 0.3 (suburbs)</td>
<td>—</td>
</tr>
<tr>
<td>Brisbane, Australia</td>
<td>0.3 (suburbs)</td>
<td>—</td>
</tr>
</tbody>
</table>

TABLE 2. Aerodynamic roughness length $z_0$ around measuring sites estimated using aerial imagery in Fig. 4 and literature.
\[ U'(t) = \sigma_U(t) \hat{U}'(t), \tag{13} \]

where \( \sigma_U(t) \) is the slowly varying standard deviation of \( \bar{U}(t) \), and \( \hat{U}'(t) \) is the reduced residual turbulent fluctuations, we cast Eq. (12) in the form:

\[ U(t) = \bar{U}(t)[1 + I_U(t) \hat{U}'(t)], \tag{14} \]

where

\[ I_U(t) = \frac{\sigma_U(t)}{\bar{U}(t)} \tag{15} \]

is the slowly varying turbulence intensity. The extraction of \( \bar{U}(t) \) from \( U(t) \) is carried out using a moving mean filter with the averaging window of 30 s (Solari et al. 2015; Burlando et al. 2017).

To investigate the abrupt changes in different scales of motion of thunderstorm winds, the proposed method in this study is applied to the large scales \( \bar{U}(t) \) and the intermediate scales of motion \( \sigma_U(t) \) (Chen and Letchford 2004), as well as the small scales of random turbulence \( \hat{U}'(t) \). The objective is to investigate the contributing factors and governing scales of motion that drive the observed changes in the instantaneous wind speed records \( U(t) \).

### 3. Results and discussion

#### a. Abrupt changes in velocity records

1) **EUROPE**

The instantaneous velocity records \( U(t) \) of several thunderstorm events from Europe are consistently divided into three or four segments using the mean (M) and standard deviation (SD) segmentation statistics (Fig. 5). Nonstationarity in the M statistics occurs only in the \( \bar{U}(t) \) and \( U(t) \) records, whereas the SD segmentation is observed in \( U'(t) \), \( \sigma_U(t) \), and \( I_U(t) \) in addition to \( \bar{U}(t) \) and \( U(t) \). The same model settings (i.e., same \( \gamma \) values) are applied to all decomposed and instantaneous velocity records for a given event. Overall, the SD segments that encompass the central peak in the \( U(t) \) records are the same length or slightly longer than the corresponding segments obtained by applying the M statistics. Only in the cases of LS of 19 April 2012 (Fig. 5e) and 25 October 2011 (Fig. 5f), the number of M and SD segments differ in such a way that the SD approach introduces an additional segment into \( U(t) \). However, the central velocity peak has the same duration in terms of both M and SD analyses. While this result suggests that the \( U(t) \) records exhibit similar M and SD changepoints, this finding is not observed in all records analyzed in the present study. Figure 5 further shows that the changepoints in \( \bar{U}(t) \) occur at the same time as in the \( U(t) \) records. Since \( \bar{U}(t) \) represents the contribution of large-scale flow structures that govern \( U(t) \); it is expected that the locations of changepoints in these two velocity records coincide in particular in the case of M statistic. The rest of the decomposed velocity records do not exhibit any abrupt changes in M if the same model settings are used as for \( U(t) \). Nevertheless, the SD approach identifies the increased wind fluctuations in \( \sigma_U(t) \) during the thunderstorm peak. The two exceptions from this observation (Figs. 5e,f) are characterized by equally pronounced wind fluctuation in the background wind before the thunderstorm segment. Figure 5e also shows that the location of changepoints in \( \sigma_U(t) \) is probably decoupled from any roughness changes due to the constant wind direction throughout the event. The wind fluctuations greatly diminished after the thunderstorm passage. In all cases, \( U'(t) \) is a stationary process without changepoints.

Another feature worth investigating in Fig. 5 is a possible relationship between changepoint locations and wind direction \( \alpha \) (i.e., changes in \( z_0 \)). In the majority of cases, it seems that these parameters are mutually independent. An example where a link between changepoints in M and SD, on the one hand, and \( \alpha(t) \) and \( z_0 \) changes, on the other hand, might exist is the LS event from 11 April 2011 (Fig. 5d). In this case, the increase in \( z_0 \) from 0.003 to 4 m, and the consequent return to 0.003 m nicely coincides with the total duration of the thunderstorm part of the record. Similarly, the SD method applied to \( \sigma_U(t) \) in Fig. 5f splits the record into two parts and the changepoint location matches the time instant of the abrupt change of \( z_0 \).

The LS thunderstorms of 5 June 2011 and 19 April 2012 (Figs. 5c,e) are characterized by relatively constant \( \alpha(t) \). These events represent rare instances when the occurrence of thunderstorm wind is not accompanied by significant changes in wind direction. However, both cases have a couple of M and SD changepoints that separate the central thunderstorm peak from the rest of the record. These events also represent the cases for which the wind direction-based segmentation procedure proposed in Gunter et al. (2017) would not differentiate between a thunderstorm and nonthunderstorm segments. The fewer cases characterized by constant wind direction in this study are expected because a shift in wind direction is one of the main signatures of a thunderstorm passage.

Figures 6 and 7 show the performances of the segmentation method applied to the events that were captured with multiple anemometers. Because such measurements are rare, in particular in Europe, the presented cases enable the analysis of spatiotemporal characteristics of these events, as well as the evaluation of the segmentation method applied to multiple velocity records from the same
event. The LI event from 1 October 2012 (Figs. 6a–c) is characterized by higher values of $I_U(t)$ at 75 m (LI5) than at 20 m AGL (LI1 and LI3), which is different from the typical turbulence profiles in ABL winds. The high values of $I_U(t)$ are caused by the sharp wind spike that occurred approximately 30 min into the time series (Figs. 6b,c). Because this velocity spike is also present in $\mathcal{U}(t)$, its existence is associated with the large-scale structures in the outflow (Burlando et al. 2017). In other words, the spike is likely not a random turbulent fluctuation. Burlando et al.
(2017) discussed that this spike in the anemometer records might represent the occurrence of a smaller-scale burst embedded into the larger-scale downburst (Sherman 1987; Hjelmfelt 1988). The absence of this feature from the LI1 record (Fig. 6a) additionally supports the hypothesis of a highly localized phenomenon.

The M and SD changepoints in the $U(t)$ and $\bar{U}(t)$ nearly overlap, except for an additional SD changepoint that appears in $U(t)$ from the LI3 anemometer (Fig. 6b). The second segment in $U(t)$ is associated with the steady increase of wind speed prior to the thunderstorm peak, and this phenomenon is sometimes discussed in terms of a gust front that precedes the thunderstorm (Droegemeier and Wilhelmson 1987; Mueller and Carbone 1987). However, the gust front feature is not observed in all velocity records. The relationship between $\alpha(t)$ and $z_0$, 

![Fig. 6](downloaded_from_http://journals.ametsoc.org/mwr/article-pdf/148/9/3747/4993433/mwrd190312.pdf by guest on 30 August 2020)
once again, shows that there seems not to be any apparent connection between those two parameters and the changepoint locations in either M or SD statistics.

Next, the LI event from 26 October 2012 (Figs. 6d–f) is particularly interesting because neither $U(t)$ nor $\overline{U}(t)$ from the LI4 anemometer (Fig. 6f) have an SD changepoint. However, the SD changepoints appear in $\sigma_U(t)$ and $I_U(t)$. This result suggests that, in rare cases, turbulent fluctuations during the thunderstorm passage might be difficult to detect (or even undetectable) in the instantaneous velocity records due to the high variability of the mean flow that governs $U(t)$ and therefore masks the changepoints in SD. Also, apart from many other events analyzed in this paper, this anemometer record is characterized by relatively constant $a(t)$ and, consequently, no change of $z_0$ occurred throughout the event.

The LI event from 4 September 2011 is the last event analyzed in this subsection (Fig. 7). The event is unusual

![Fig. 7. As in Fig. 5, but for a thunderstorm event from Livorno captured with five anemometers.](http://journals.ametsoc.org/mwr/article-pdf/148/9/3747/4993433/mwrd190312.pdf)
because while the $U(t)$ records from all anemometers are visually very similar, the changepoint method shows that the statistical properties are different among the records. Although the same model settings are used in all velocity records in Fig. 7, the number and location of changepoints are not the same. Therefore, the analyzed thunderstorm outflow was highly localized because statistical properties are different from one location in the outflow to another. Similar to Fig. 6, this event was also characterized by the constant value of $z_0$ throughout the record.

2) THE UNITED STATES

Figure 8 shows the abrupt changes in several velocity records from the United States. In comparison to the events from Europe that are usually associated with single- and multicell thunderstorm clouds (Burlando et al. 2017, 2018), the PE, SE, and LU (21 May 2008 and 4 June 2009) events were supercells with rear flank downdrafts, misovortices, or mesoscale convective systems such as bow echoes and mesoscale convective complexes (Holmes et al. 2008; Lombardo et al. 2014; Gunter and Schroeder 2015; Gunter et al. 2017). Further, the comparison of PE and SY events against the rest of LU records in Fig. 8 is also appealing due to the discrepancies in the length of the velocity records—the former is about 6 times longer. Despite this difference in length, the dominant thunderstorm velocity peaks are identified in all $U(t)$ records using the M method. In some cases (e.g., Figs. 8a,d,e,f), wind fluctuations extend beyond the first peak, and the SD method accordingly separates the larger portion from the entire velocity record. The M changepoints in $\bar{U}(t)$ and $U(t)$ always occur at the same position in the records. In the PE event, a secondary circulation on the leading edge of the outflow contributed to the multiple peaks in the velocity record (Gunter et al. 2017). This result, once again, demonstrates that the abrupt changes in the mean velocity are driven by the large-scale flow structures in thunderstorm winds. While the interpretation of changepoints in the rest of decomposed records in Fig. 8 is more challenging than the explanation of changepoints in $\bar{U}(t)$, we notice the inconsistency in the location of changepoints in $I_U(t)$. While in some cases $I_U(t)$ is a stationary process in terms of SD (Figs. 8b,f), in other cases, there are sharp changes of SD in the $I_U(t)$ records (Figs. 8c,d,e). These changepoints cannot be associated with the abrupt changes of $z_0$ for different values of $\alpha$ because the sites are characterized by the homogeneous and isotropic roughness. Because the anemometer heights were also comparable, it is not clear if this factor could be the governing contributor for the location and number of changepoints in this case.

Last, we briefly analyze the LU event from 4 June 2002 that was measured with multiple anemometers. Figure 9 shows the changepoints in 6 out of 12 velocity records from this event. All measurements were constrained below the heights of 15 m, where the influence of surface roughness is deemed to be significant. The $U(t)$ records show a high degree of similarity in terms of the number and locations of changepoints. In all cases, these 30-min records are characterized by two M changepoints (the thunderstorm peak) and a maximum of one SD changepoint. The SD changepoints, in all cases, separate the thunderstorm wind segment from the winds after the thunderstorm. The changepoint locations are consistent in all $U(t)$ records and the absence of changepoints is noticed in the records of $I_U(t)$ and $U'(t)$.

3) AUSTRALIA

The velocity records from Australia are only analyzed in terms of $U(t)$ due to the low sampling frequency of obtained data (Table 1). All $U(t)$ records from BR are shown in Fig. 10. The similarity between time histories of $U(t)$ at 58 and 104 m AGL (Figs. 10a,b) is higher regarding the location and number of changepoints than between either of the two records and 10-m data from the nearby mast (Fig. 10c). However, because the duration of these velocity records is also different, the similarities are discussed only in terms of the central 10-min velocity peak. The analysis of numerous records whose duration is $\geq$ 60-min has shown that 10-min records are not sufficiently long to accurately represent the entire transient nature of thunderstorm winds and their separation from the background winds before and after the event. We again observe that the SD portion is longer than the M segment that envelops the central velocity peak.

b. Spectral analysis

The turbulence spectra obtained from all $U'(t)$ records with $f_s \geq 1$ Hz are shown in Fig. 11. The LU thunderstorm of 4 June 2002 is the only U.S. event that entirely entered into this analysis due to the anemometer response issues reported in the other U.S. events (Lombardo et al. 2014; Gunter et al. 2017). The power spectral density (PSD) analysis of the Lubbock events from Lombardo et al. (2014) is limited to the frequency range 0.04–1 Hz due to the internal anemometer sampling procedures and anemometer response issues. Also, mechanical filtering and second-order anemometer response characteristics resulted in unreliable energy spectra between 0.4 and 5.4 Hz for the PE and SY events (Gunter et al. 2017). Propeller-vane anemometers mechanically filter wind speed in proportion to the anemometer’s distance constant due to the propeller inertia. The distance constant of an anemometer is the length of airflow past the sensor required to cause it to respond to 63.2% of the increasing step function change in wind speed.
Gunter et al. (2017) and Schroeder and Smith (2003) showed that the distance constant of 2.7 m diminishes a wind gust with a wavelength of 30 m by 13%, and a wind gust of 17 m wavelength by approximately 71%. That is, the anemometer mechanically filters the amplitudes of short wavelength gusts.

These unreliable segments of PSDs are not included in Fig. 11. Here, the PSDs are evaluated using the Welch (1967) method of modified periodograms with a Hamming window of $2^6$ velocity readings. A 50% overlap between adjacent windows was applied. While all thunderstorm records from Europe are characterized by a $-5/3$ slope of PSD in the inertial subrange (Fig. 11), the steeper slope of PSD is found in the U.S. events. This deviation from the typical ABL spectra was also reported in Holmes et al. (2008) and Lombardo et al. (2014). The larger slope in the
events from Lombardo et al. (2014) is found at the frequencies higher than approximately 0.375 Hz. Also, there seems to be a well-established peak centered around 0.8 Hz in all $\bar{U}(t)$ records from the LU event of 4 June 2002 (not explicitly shown). The peak origin and the deviation of PSD slope from the reference slope in the U.S. data are not fully understood at this point. One possible explanation often presented in the literature is the applicability of Eq. (12) and the 30-s averaging window in separating low- from high-frequency fluctuations (Holmes et al. 2008; Lombardo et al. 2014). This issue was thoroughly addressed in Solari et al. (2015). The peak and the steeper slope could also be caused by an anemometer issue or insufficient data quality that stayed unknown to the authors.

The PSD flattening for the European events in the highest frequency domain was also reported in Burlando
et al. (2017) and a likely cause is spectral aliasing and anemometer noise. For example, this feature is not observed in the PSDs from HY (Finland) and LS (Italy) events (not explicitly shown Fig. 11). However, it should also be noted that the spectra and other results discussed herein are not necessarily the “typical” representative of thunderstorm velocity records in these regions due to the limited number of cases that are considered in this research.

The segmentation methodology also enables the PSD analysis of specific segments in the \( U(t) \) records in which the M and SD changepoints occurred at the same or similar times (e.g., HY in Fig. 5a, LS of 5 October 2011 in Fig. 5c, LI3 of 26 October 2012 in Fig. 6e). Figure 12 shows the results of this analysis on the example of HY data. Since the PSDs in Fig. 12 are not normalized; their values are proportional to the velocity fluctuations in the represented segment. The PSD of thunderstorm peaks (black line) is associated with the highest fluctuations and more turbulent energy than the rest of the velocity record. At a given frequency, the separation between PSD of the thunderstorm peak and the PSDs of the winds before (dark gray) and after (light gray) the peak is proportional to the fluctuation differences in the corresponding velocity segments. Analyzing multiple anemometer records, we noticed that there is no clear distinction between fluctuations in the winds before and after the thunderstorm from one event to another (not shown in Fig. 12). In some cases, the velocity fluctuations before the thunderstorm peak are more pronounced than in the winds after the peak, and sometimes the reverse. These results cannot be a priori attributed to surface roughness changes due to the constant wind direction in some of the investigated cases.

c. Model applications

The proposed segmentation method is objective and accurate at separating the so-called intermediate wind events from the thunderstorm winds (De Gaetano et al. 2014). The intermediate winds are defined as a stationary and non-Gaussian wind, whereas the thunderstorm winds are nonstationary and non-Gaussian. As discussed in the Introduction, the De Gaetano’s methodology distinguishes thunderstorm and intermediate winds from the ABL winds using the 10-min and 1-h values of various statistical parameters averaged over 1-s and 1-min intervals, respectively. However, further separation between the thunderstorm and intermediate winds is afterward manual and depends on visual examination of the velocity records. Figure 13 shows two intermediate wind records, both from La Spezia, Italy (De Gaetano et al. 2014). These records are similar to thunderstorm winds in terms of gust factors and higher-order statistics (De Gaetano et al. 2014; Jubayer et al. 2019), but visually the time histories are profoundly different from the thunderstorm winds that occurred in La Spezia (Figs. 5c–e). The intermediate winds—while evidently

FIG. 10. Changepoints in the M and SD statistics in the \( U(t) \) records from the Brisbane event in Australia. (a)–(c) Anemometers BR1–3, respectively.

FIG. 11. Power spectral density (PSD) functions of all \( U(t) \) records with \( f_s \lesssim 1 \) Hz from the United States (gray) and Europe (black). The black dashed line is the \(-5/3\) slope.
different from thunderstorms—remained undistinguishable from the thunderstorm winds in the De Gaetano’s wind classification method. However, the application of the proposed segmentation methodology to the intermediate winds in Fig. 13 shows a complete absence of changepoints in both M and SD statistics. Therefore, the absence of changepoints in these types of winds can be implemented to the De Gaetano’s procedure to separate the thunderstorm from intermediate winds. The methodology is objective and certainly faster than the subjective visual interpretation that is currently used. It is equally important to emphasize that the penalty constant $g$ is set to the values of 20 000 and 10 000 for the M and SD statistics, respectively, which were also the values used to detect multiple changepoints in the thunderstorm wind records from the same anemometer stations.

We also compare our methodology against the segmentation method proposed in Lombardo et al. (2014). Both methods almost identically identified the beginning of the thunderstorm segment (Fig. 14), but our method extends the duration of the first peak in terms of both M and SD statistics in comparison to the Lombardo et al. (2014) methodology. Also, the existence of the second thunderstorm peak is identified in the M approach, whereas this segment is absent in Lombardo et al. (2014) and SD statistics. While the Lombardo’s method inspects the stationarity of data by performing the reverse arrangement and run tests on data segments, the current methodology recursively minimizes the residual error of segments in terms of the given statistics (M or SD), as described in section 2b. Without an attempt to determine which methodology is “better” at detecting changepoints, we notice that the method...
The number of detected changepoints depends only on $\gamma$. Fewer or more changepoints can be detected by increasing or decreasing $\gamma$, respectively. More changepoints being detected in Fig. 14 also indicates that the proposed method can identify finer-scale features in the velocity records. Last, the proposed method in this article is tested for a broader range of features in the velocity records. The introduced model for the detection of changepoints in thunderstorm wind records is robust because the number and location of changepoints in a velocity record are unchanged for a wide range of frequency records. The second objective was imposed to make the model robust and consistent for a wide range of thunderstorm winds.

Figure 15 shows that the proper segmentation of the GE event (30 September 2012; Fig. 5b) is obtained for $\gamma_M \approx 6000$ (subscript “M” stands for the mean cost function) and the further increase of $\gamma_M$ does not introduce additional changepoints. However, the model is sensitive to $\gamma_M < 5000$ and a different number of changepoints is found depending on the exact value of this parameter. This relationship is more subtle in the case of high-frequency data, such as the LU event from 4 June 2009 (Fig. 15), because these measurements can reveal more finescale structures in the flow. In this event, the number of changepoints decreases from 65 to 8 for the values of $\gamma_M$ ranging between 20000 and 1000.

The following analysis further investigates the relationship between $\gamma_M$ and $f_s$ in Figs. 16 and 17. Here, decimation is used to reduce artificially $f_s$ of measured data. The decimation method employs a low-pass Chebyshev type I infinite impulse response filter of order 8 (Parks and Burrus 1987) to reduce the signal bandwidth and guard against aliasing (Jackson 1996). This analysis demonstrates that, in order to preserve the same number and location of changepoints in the decimated velocity records, $\gamma_M$ needs to decrease with reducing $f_s$. This relationship between $\gamma_M$ and $f_s$ for the M cost function is linear and in the form:

$$\gamma_M = \gamma_{ref} f_s,$$

where

$$\gamma_{ref} = \gamma_M|_{f_s=1Hz}.$$
e. Transient characteristics of thunderstorm velocity records

The duration of the thunderstorm wind peak and the ramp-up times are shown in Fig. 19. Here, the duration of thunderstorm velocity peak $\Delta T_p$ is defined as the length of M or SD segments that contain the dominant peak in the velocity record. The ramp-up time $\Delta T_{up}$ is the interval between the first changepoint before the peak velocity and the time of the peak. This definition of ramp-up time is similar to the rise time often used in the analyses of turbulent velocity peaks in ABL winds (Hu et al. 2018).

The peak duration seems to be similar between the U.S. and European events (Fig. 19). In many cases, the SD segments of the thunderstorm peak are longer than the associated M segments. While this difference in segments duration is not always pronounced, Fig. 19a demonstrates that the fluctuations often precede and proceed the abrupt changes of M. The duration of thunderstorm velocity peak is 2–5 min in approximately 60% of the analyzed records using the M cost function (Fig. 19b), while the peak duration is 5–10 min in approximately 60% of the same records that are partitioned using the SD approach (Fig. 19c). Besides, there is not a velocity record in which the SD segment around the thunderstorm peak is shorter than 2 min. The velocity records that are characterized by a single changepoint were not considered in this analysis because the peak duration cannot be evaluated in those cases.

The ramp-up times are below 2 min in approximately 60% of the analyzed velocity records using the M statistics (Fig. 19e). Similar to the duration times, the SD method isolates the longer ramp-up times than the M cost function (less than 10% of the records are characterized by the ramp-up time below 1 min). Therefore, the thunderstorm winds seem to be more transient in terms of the abrupt changes of M than SD. Some events, such as the LU 4 June 2002, are characterized by a single changepoint in the SD statistic (Figs. 19a,d). Despite the low sampling frequency of the BR3 event from Australia and the WA event from the United States, these cases are still characterized by the similar peak duration and ramp-up times as the other events. This result shows that
the main transient features of thunderstorm records are noticeable even in the low-frequency data. Moreover, this analysis demonstrates that the proposed segmentation methodology is also applicable to low-frequency data.

Another parameter of importance in the analysis of thunderstorm winds is the velocity speedup during the thunderstorm wind episode. We quantify this speedup by defining a ratio of the maximum value of the slowly varying mean velocity during the thunderstorm segment $\hat{U}_p$ and mean wind speed prior to the thunderstorm $U_b$ (the subscript “$b$” stands for background):

$$R_m = \frac{\hat{U}_p}{U_b}. \quad (18)$$

Similarly, we also define a ratio between the standard deviation of instantaneous velocity during the thunderstorm peak $\sigma_p$ and that during the background wind prior to thunderstorm $\sigma_b$ as

$$R_f = \frac{\sigma_p}{\sigma_b}. \quad (19)$$

The above ratios are only defined for the velocity records that have a clear distinction between the thunderstorm peak and the rest of the time series (i.e., more than one changepoint). The background segment is the interval between the beginning of the record and the first changepoint. For instance, the LU event of 4 June 2002 did not enter into this analysis due to the short duration of the velocity record and unclear existence of the background wind prior to the thunderstorm segment. The ratios in Eqs. (18) and (19) are similar to those defined by Solari et al. (2015) in their study of thunderstorm characteristics that are of importance in wind engineering. However, the ratios proposed herein directly quantify the abrupt changes in $M$ and SD between the thunderstorm segment and the non-thunderstorm winds prior to the peak.

The maximum velocity during the thunderstorm peak in the events from Europe is typically 2–4 times larger than the mean wind speed before the event (Fig. 20). However, the HY event from Finland and the GE event from Italy are the exceptions because the ratio is much higher. The ratio $R_f$ is smaller than $R_m$ in the majority of analyzed records and, on average, $R_m \approx 1.65R_f$. Therefore,
the abrupt changes associated with the largest scales of motion (represented through the mean and moving mean) are more pronounced than the changes in the smaller scales of motion (represented through velocity fluctuations). This result independently confirms (i.e., not through the length of isolated segments that was discussed previously) that the abrupt changes of M are more pronounced than the SD changes during the thunderstorm wind episodes. The U.S. events are characterized by slightly larger $R_m$ than European records. We did not observe any apparent relationship between $R_m$ or $R_f$ and $z_0$.

The reported ratios are also crucial for the proper physical reconstruction of thunderstorm winds in wind tunnels. Of particular importance are the wind simulators capable of simultaneously producing thunderstorm downbursts and background ABL winds (Romanic et al. 2019). Namely, the ratios impose the relative strength that the two simulated flows need to have in terms of mean velocity and turbulent fluctuations in order to satisfy the real atmospheric conditions. In most cases, the relative strength of the flows needs to be such that the peak velocity during the thunderstorm downburst is approximately 2–4 times higher than the mean wind speed in the background winds prior to the downburst.

4. Conclusions

This study investigated the transient characteristics of thunderstorm winds by using an objective method for the detection of changepoints in velocity records. The method seeks changepoints in anemometer records based on the statistical properties of different segments of time series (Lavielle 2005; Killick et al. 2012). The investigated wind statistics were the mean (M) and standard deviation (SD). We considered 41 velocity records from 19 thunderstorm wind events from Europe (Järvi et al. 2007; Solari et al. 2012), the United States (Fujita 1985; Orwig and Schroeder 2007; Lombardo et al. 2014; Gunter and Schroeder 2015), and Australia (Sherman 1987). The changepoint analysis was applied not only to the instantaneous velocity records but also to the decomposed time series in the form of slowly varying mean, moving standard deviation and moving turbulence intensity, as well as the residual turbulent fluctuations. The main conclusions of this study are briefly summarized below.

- The proposed methodology separated the central thunderstorm wind peak from the rest of the velocity records. The M and SD changepoints in the instantaneous velocity records coincide with the location of changepoints in the records of slowly varying moving mean, moving standard deviation and moving turbulence intensity, as well as the residual turbulent fluctuations. The main conclusions of this study are briefly summarized below.
FIG. 19. (a) The duration of thunderstorm peak $\Delta T_p$ obtained using the M and SD cost functions. (b), (c) Associated histograms of $\Delta T_p$. (d)–(f) Similarly to the peak duration, these are the duration and histograms of the ramp-up times $\Delta T_{up}$.
This systematic observation suggests that the changes in velocity fluctuations often precede and proceed the abrupt changes in the mean flow.

This study did not find any methodical relationship between the time occurrence of changepoints in either M or SD statistics and the surface roughness changes associated with the shifts of wind direction during the thunderstorm wind episodes. It seems that the thunderstorm outflows might be too localized and transient to establish the flow equilibrium with the underlying surface. The absence of systematically detected changepoints in the records of slowly varying turbulence intensity and standard deviation for continually changing wind direction during some thunderstorm episodes partially supports the conclusion that the time evolution of thunderstorm winds is rapid to the point the flow field does not reach equilibrium with the terrain roughness. While similar results were also reported in Burlando et al. (2017), more research is needed on the dynamics and evolution of thunderstorm winds over the zones of abrupt roughness changes.

The study further showed that the maximum velocity during the thunderstorm peak in the events from Europe is usually between 2 and 4 times larger than the mean wind speed before the event. This ratio seems to be higher for U.S. events, but the analyzed sample was too small to conclude this result firmly. Also, the U.S. events were more organized thunderstorm systems than the isolated thunderstorms in Europe, and this difference in thunderstorm morphology is probably an important factor that governed the above ratios. The duration of the thunderstorm velocity peak was 2–5 min in approximately 60% of the records segmented using the M cost function and 5–10 min long when the records were partitioned using the SD statistic.

Also, the ramp-up times of the mean flow were below 2 min in approximately 50% of the analyzed velocity records. Here, the ramp-up time was defined as the time interval between the velocity peak in the slowly varying velocity record and the first changepoint before the velocity peak. The ramp-up time was longer when the records were analyzed using the SD cost function.

Last, this study provides the proper configuration of the proposed method for the detection of changepoints in velocity records. The values of threshold parameter $\gamma$ were derived for both M and SD cost functions, as well as different record durations. The functional dependency between $\gamma$ and the anemometer sampling frequency was also determined for the M and SD statistics. Practical applications of the method were demonstrated by objectively separating intermediate wind events from thunderstorm downbursts in the De Gaetano et al. (2014) methodology of wind type classification. Currently, their methodology differentiates between these two wind systems through visual observation of the records (i.e., an expert’s opinion). Also, the performances of the proposed segmentation method were compared against the results of two similar methods employed by Lombardo et al. (2014).

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**Data availability statement.** Data analyzed in this study were a reanalysis of existing data, which are available at locations cited in the reference section.

**REFERENCES**


